

特定領域研究
「スーパークリーン物質で実現する
新しい量子相の物理」
A03,A04班 合同研究会

19th April, '09
Hotel Hakone Powell

Ferromagnetic Superconductor UCoGe *Co NMR/NQR*

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Ferromagnetic Superconductor

Ferromagnetism & Superconductivity

(with a spin-singlet pairing)

Mutually exclusive ?!

Coexistence of ferromagnetism and superconductivity has been known.

Examples

$(Ce_{1-x}Gd_x)Ru_2$ ('58)

$x=0.12$

$T_{SC} \sim 5\text{ K}$

$T_{Curie} \sim 4\text{ K}$

$RuSr_2GdCu_2O_8$ ('96)

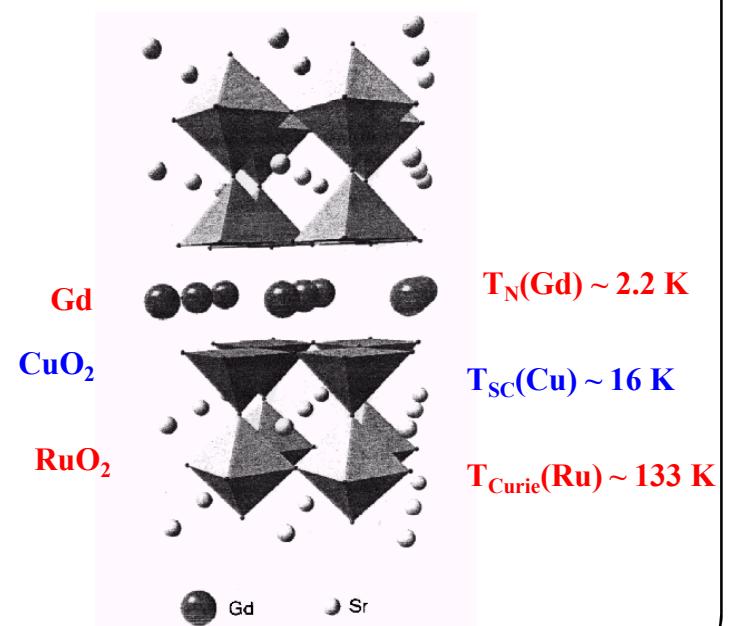
magnetic ordering

$T_{Curie}(Ru) \sim 133\text{ K},$

$T_N(Gd) \sim 2.2\text{ K}$

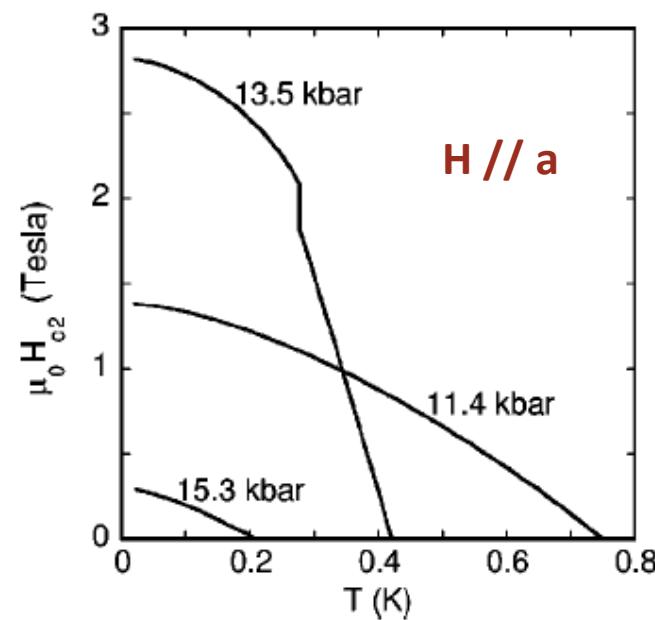
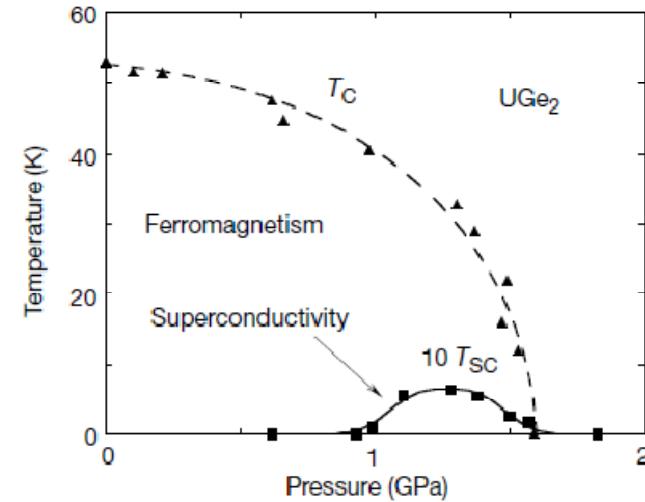
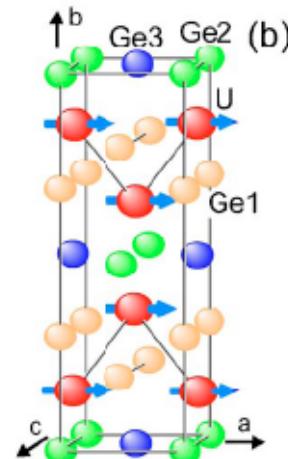
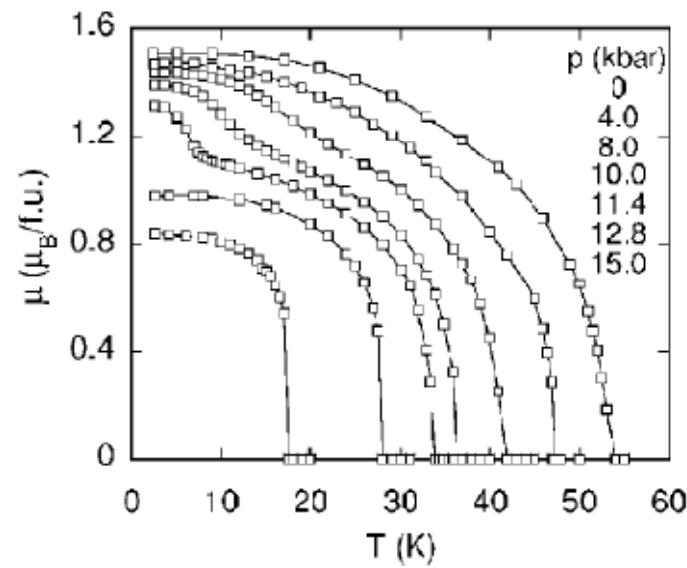
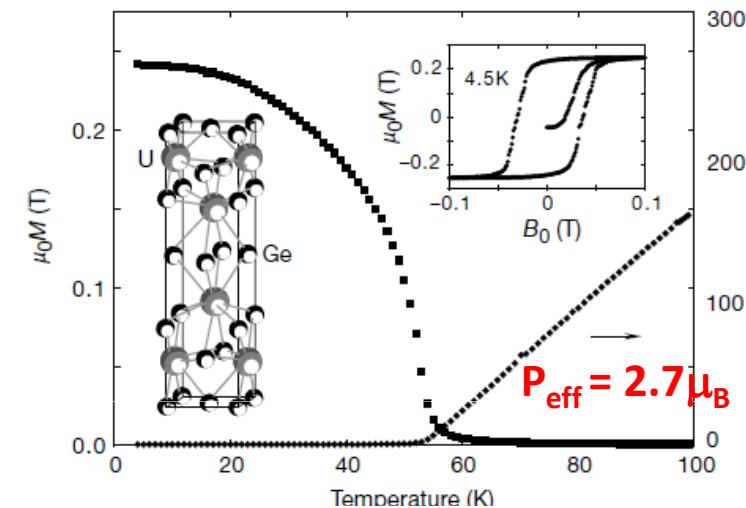
SC transition

$T_S(Cu) \sim 16\text{ K},$



Break Through (S.S. Saxena *et al.* Nature '00)

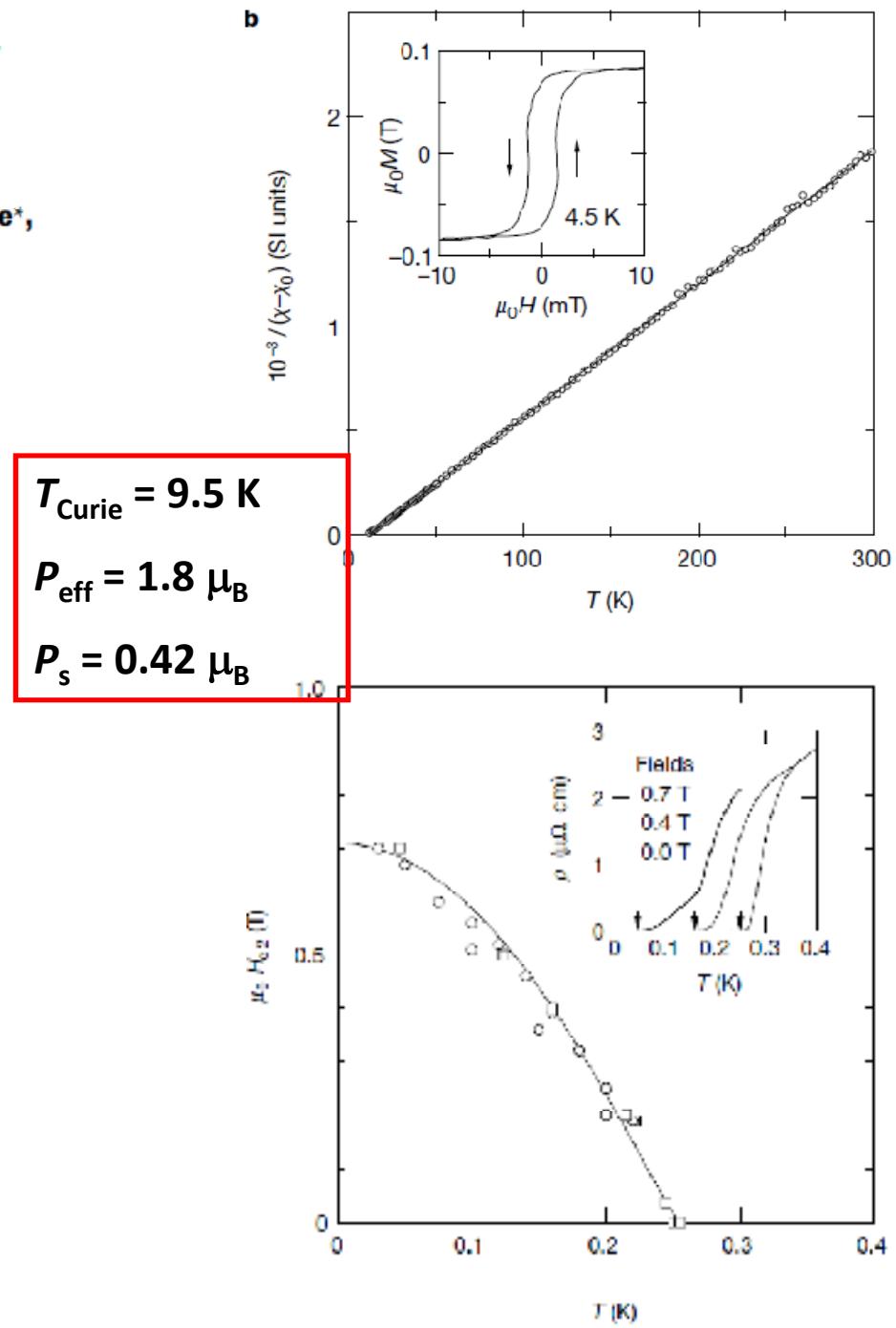
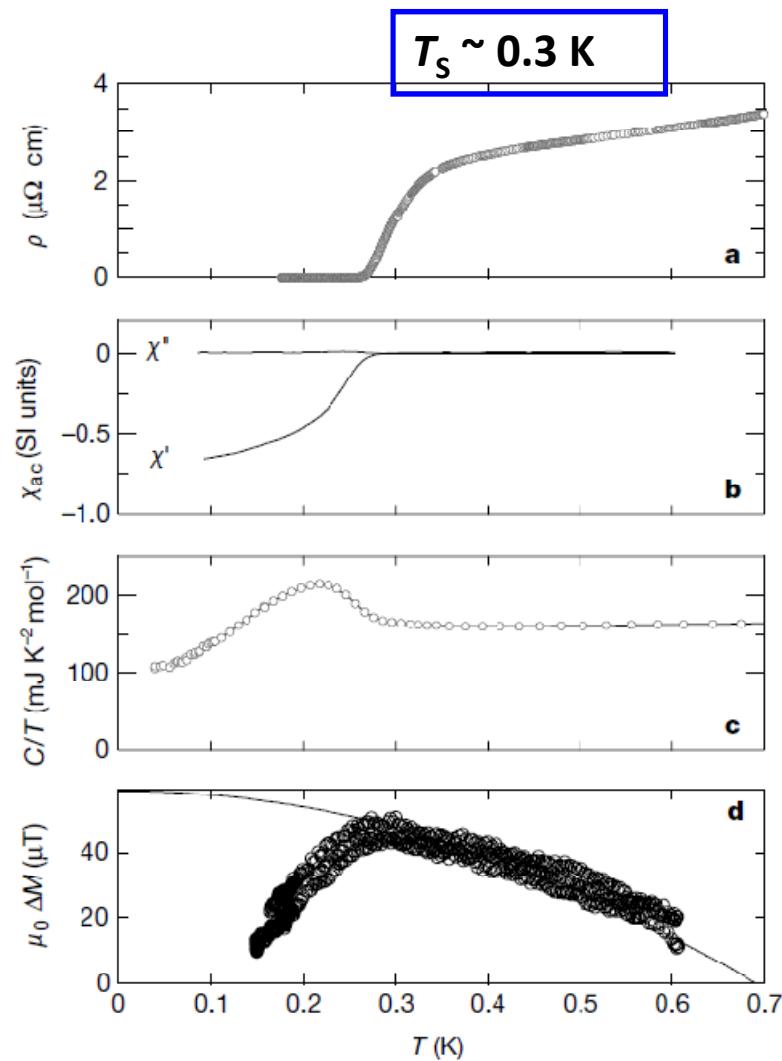
UGe₂



Sheikin *et al.*

Coexistence of superconductivity and ferromagnetism in URhGe

Dai Aoki*, Andrew Huxley*, Eric Ressouche*, Daniel Braithwaite*, Jacques Flouquet*, Jean-Pascal Brison†, Elsa Lhotel† & Carlev Paulsen†





Superconductivity on the Border of Weak Itinerant Ferromagnetism in UCoGe

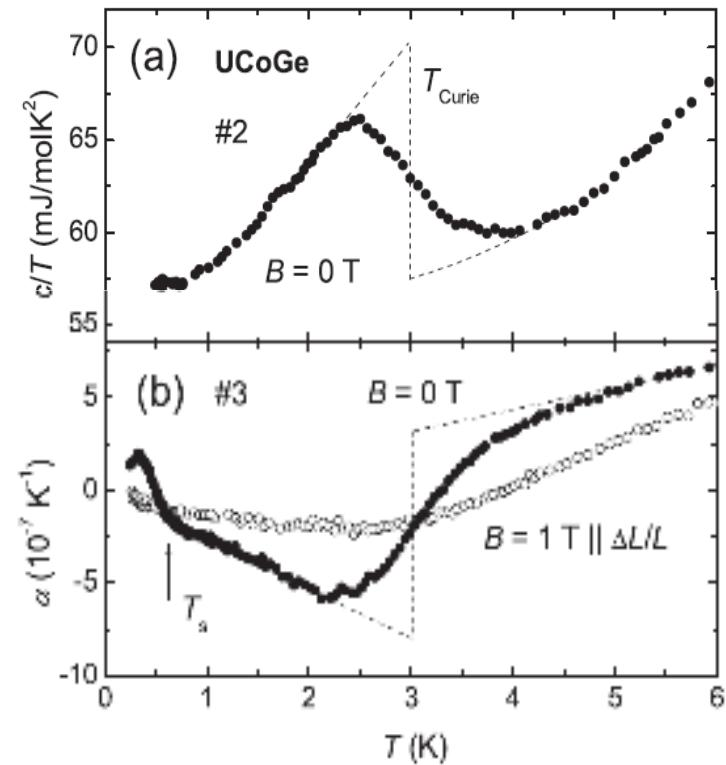
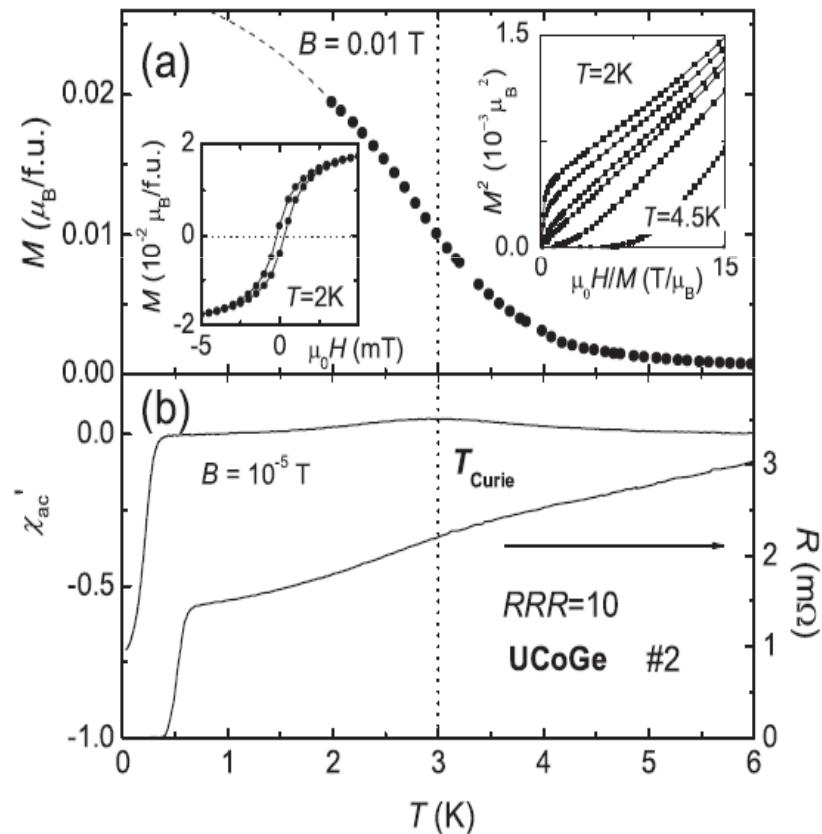
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$$T_C = 3 \text{ K}$$

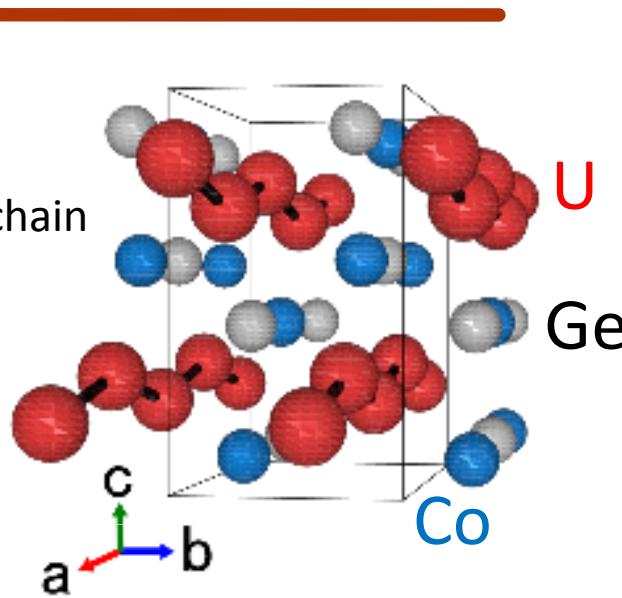
$$T_S = 0.8 \text{ K} \text{ (best sample)}$$

$$m_0 = 0.03 \mu_B : \text{weak itinerant FM}$$

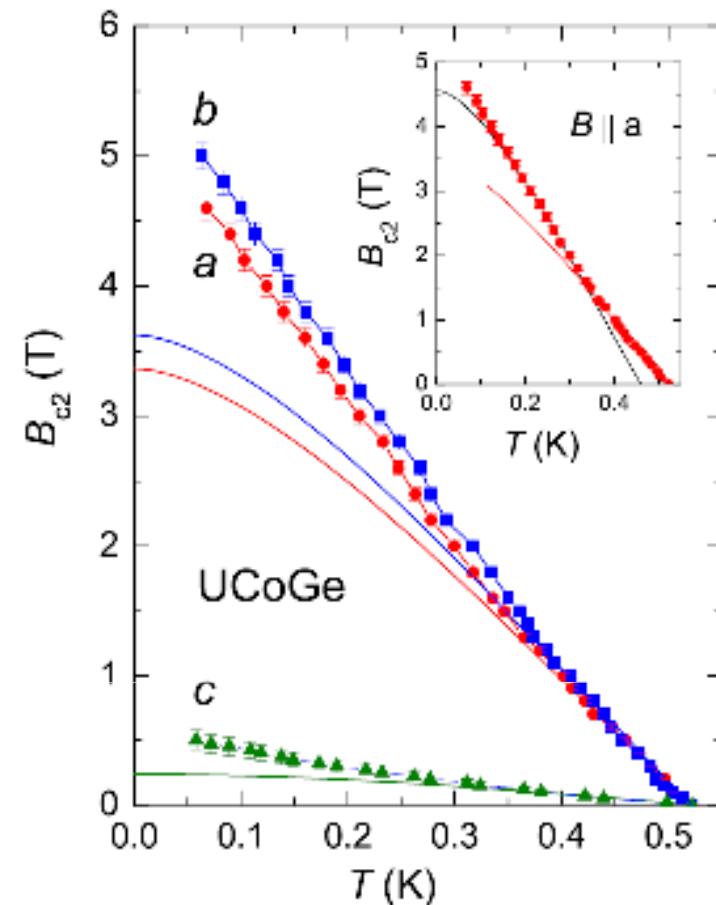
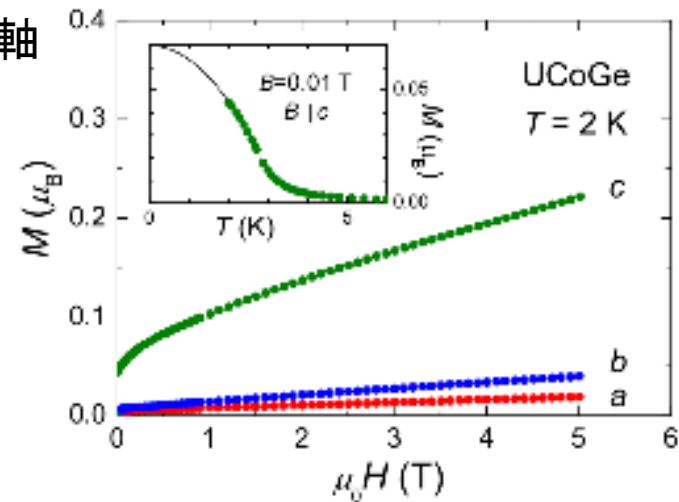
Introduction 2



- Othorhombic
- **U** forms zigzag chain



Easy axis: C軸



N. T. Huy, et al., Phys. Rev. Lett. 100 077002 (2008)

Large Upper Critical Field:
Greater than Pauli-Limit field (that is an expected H_{c2} in the spin-singlet SC)

$$(B_{c2}^{\text{Pauli}}(0K) = 1.83T_{SC} \sim 1\text{ T})$$

Spin-triplet superconductivity?

Motivation

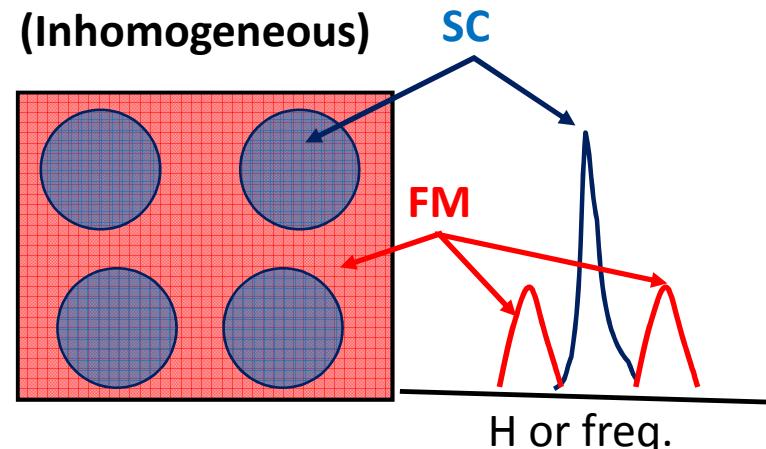
- By using Co-nuclear quadrupole resonance_(NQR), we have investigated magnetic and SC properties in UCoGe.
 - Whether SC and FM states coexist microscopically or not?
⇒ Microscopic measurements are crucial.

Sample: Provided by

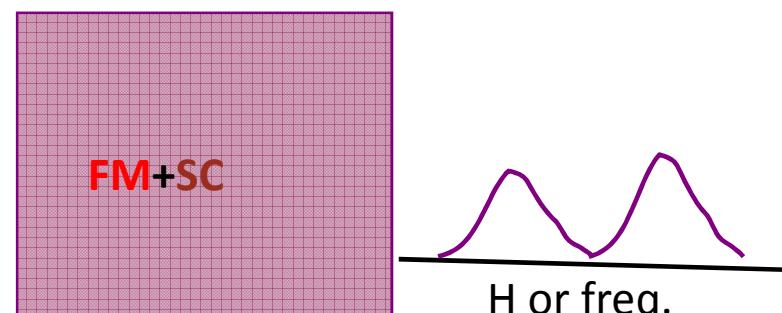
佐藤憲昭, 出口和彦 @名古屋大学理学研究科
佐藤伊佐務 @東北大学金属材料研究所

$$T_{\text{curie}} > T_{\text{SC}} > T$$

Compete
(Inhomogeneous)



Coexistence
(homogeneous)



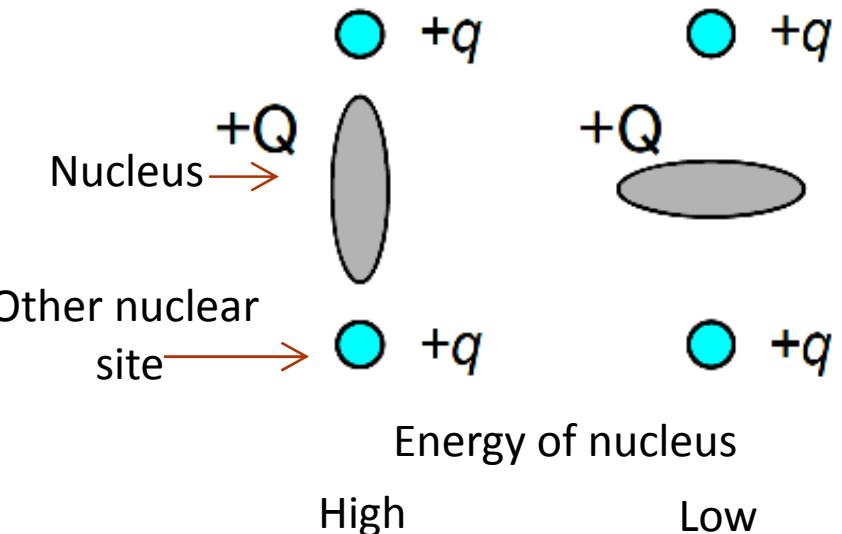
Nuclear Quadrupole Resonance (NQR)

When nuclear Spin ($I \geq 1$)
Electric quadrupole moment (eQ) is present



Energy Levels are split.

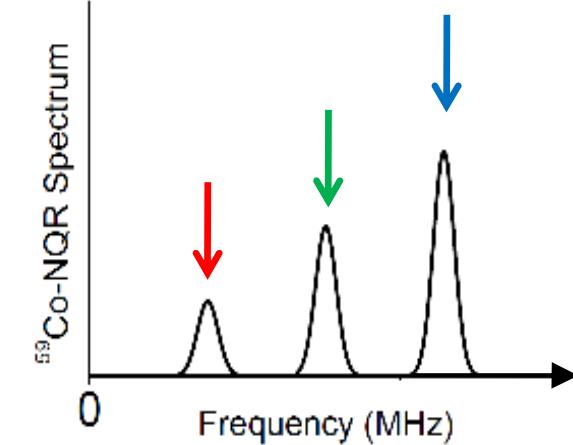
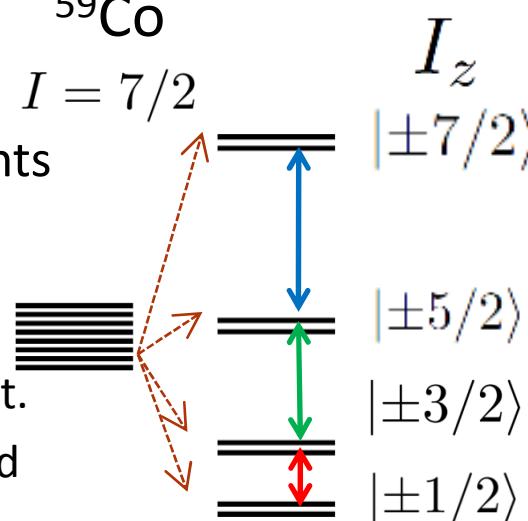
When RF fields corresponding to the energy splitting are applied, resonance is observed.



Advantage of NQR Measurements

- Zero external fields
⇒ SC is not suppressed
- Signal from electric quadrupole Int.
⇒ Microscopic information around nuclear site

^{59}Co
 $I = 7/2$

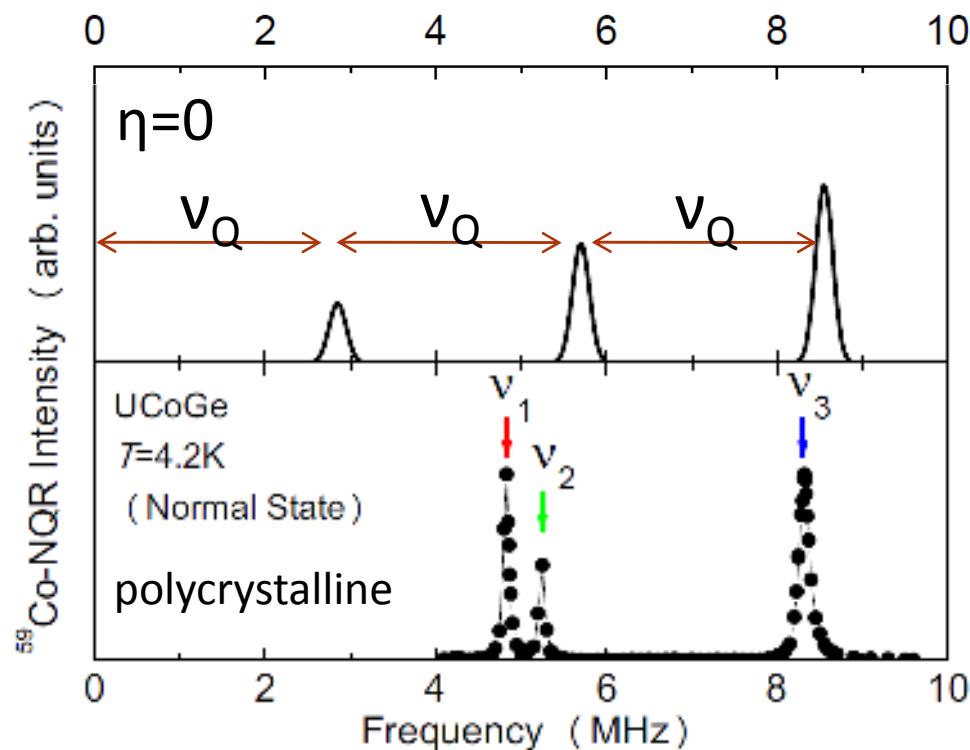


^{59}Co -NQR spectrum 1

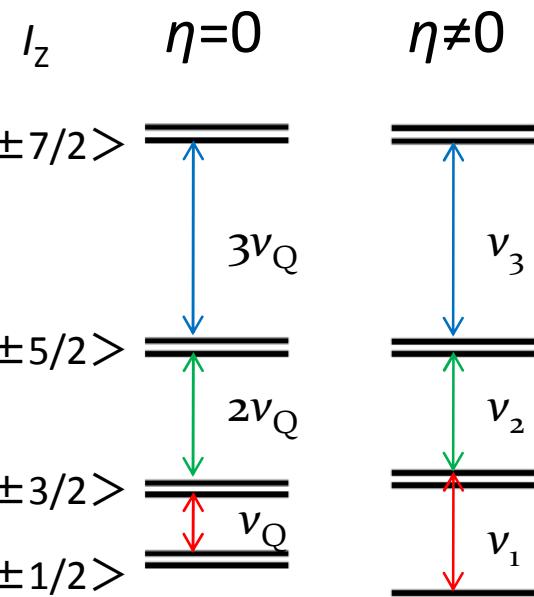
$$\hat{H}_Q = \frac{\nu_Q}{6} \left\{ \left(3\hat{I}_z^2 - \hat{I}^2 \right) + \frac{1}{2}\eta \left(\hat{I}_+^2 + \hat{I}_-^2 \right) \right\}$$

Asymmetric parameter

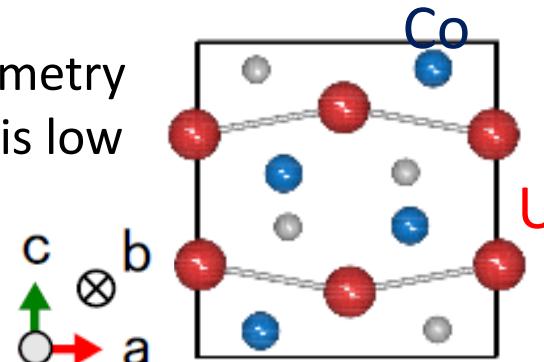
$$\eta = \frac{V_{XX} - V_{YY}}{V_{ZZ}}$$



^{59}Co (nuclear spin $I = 7/2$)



Local symmetry at Co site is low
⇒ $\eta \neq 0$



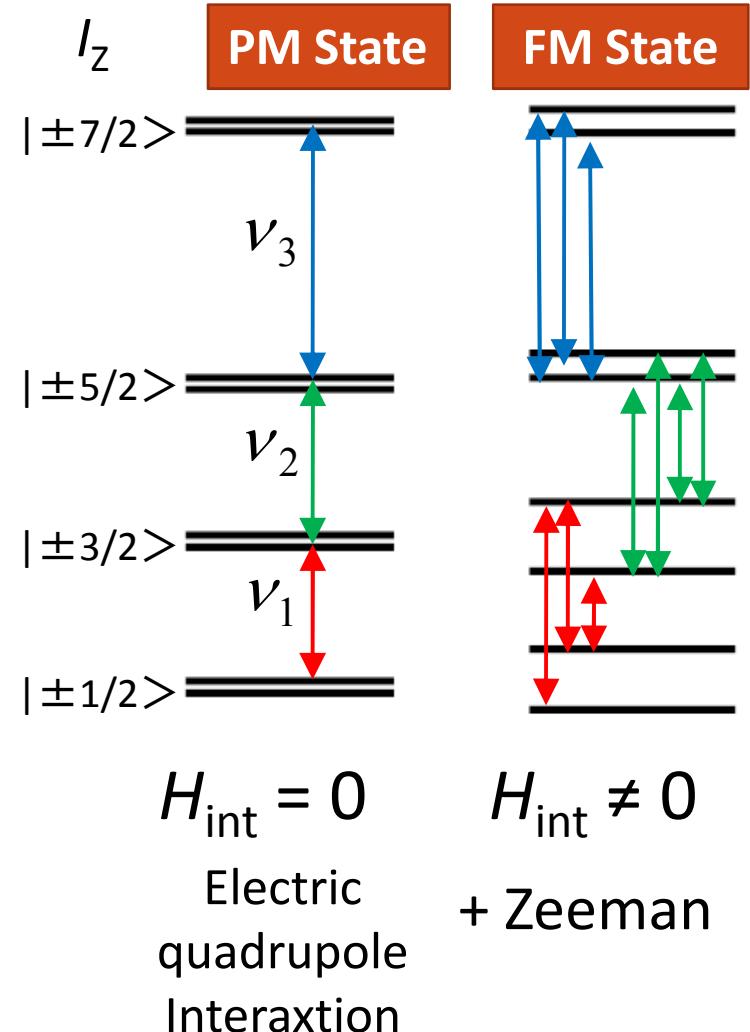
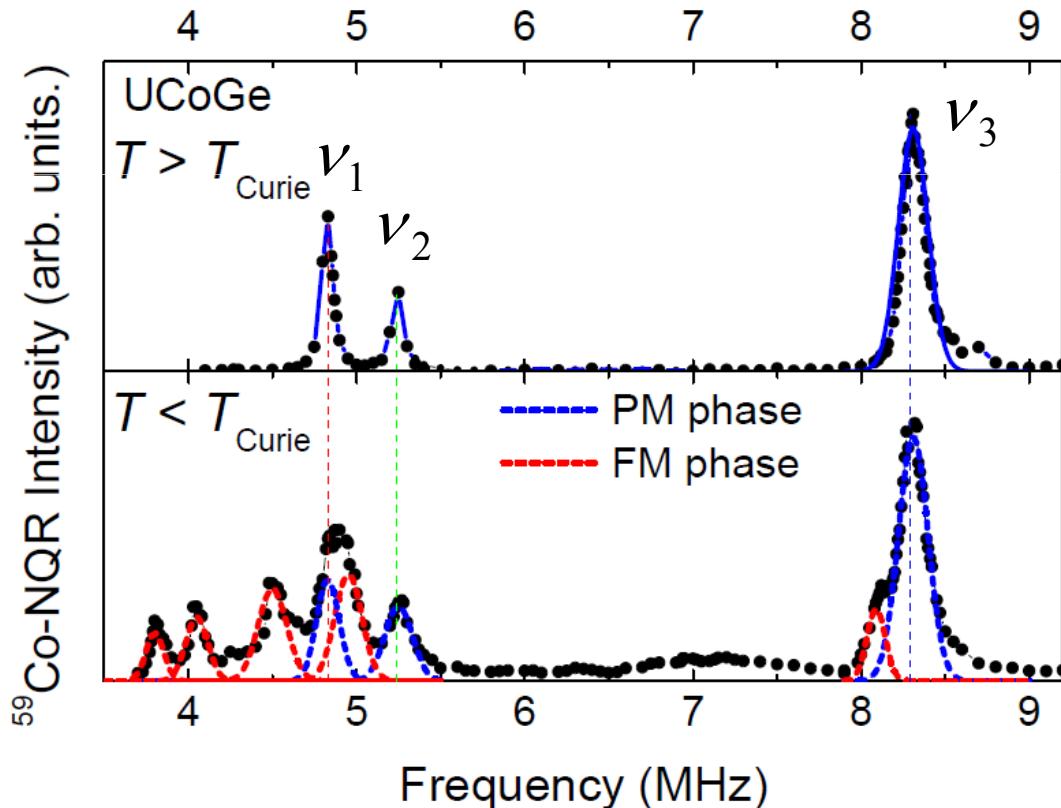
$\nu_Q = 2.85\text{MHz}$, $\eta = 0.52$

^{59}Co -NQR spectrum 2

FM state

$$\mathcal{H} = \mathcal{H}_Q - \underline{\gamma \hbar \mathbf{I} \cdot \mathbf{H}_0}$$

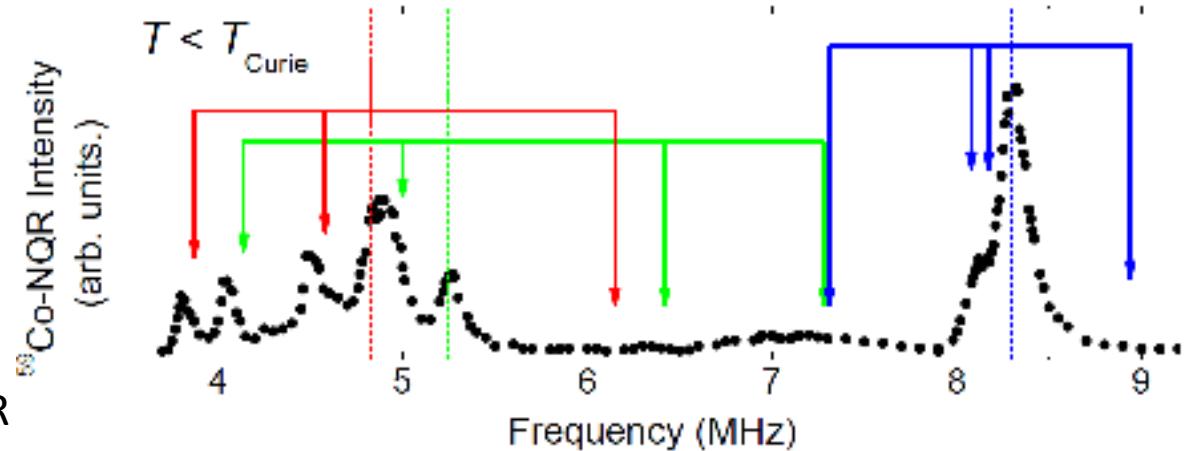
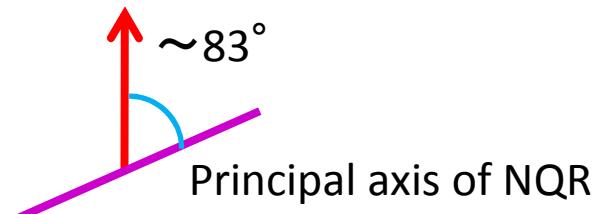
Zeeman Interaction by internal field



Internal Field

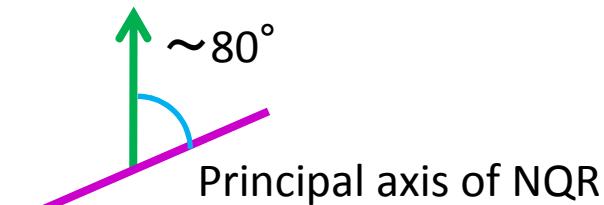
NQR spectrum

H_{int} at ${}^{59}\text{Co}$ (~ 900 Oe)

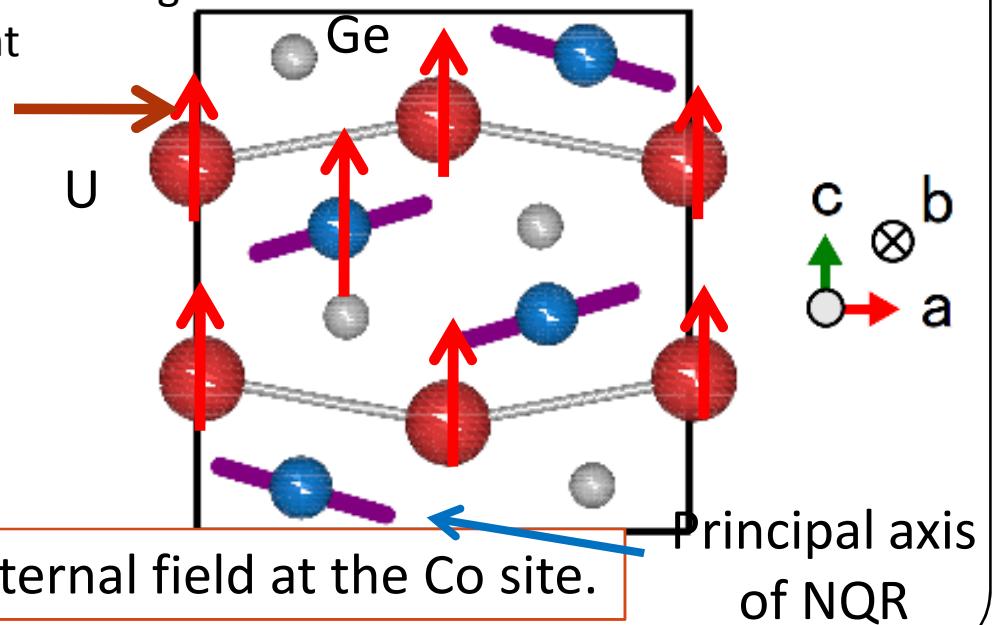


Band Cal. (by. H. Harima)

c 軸

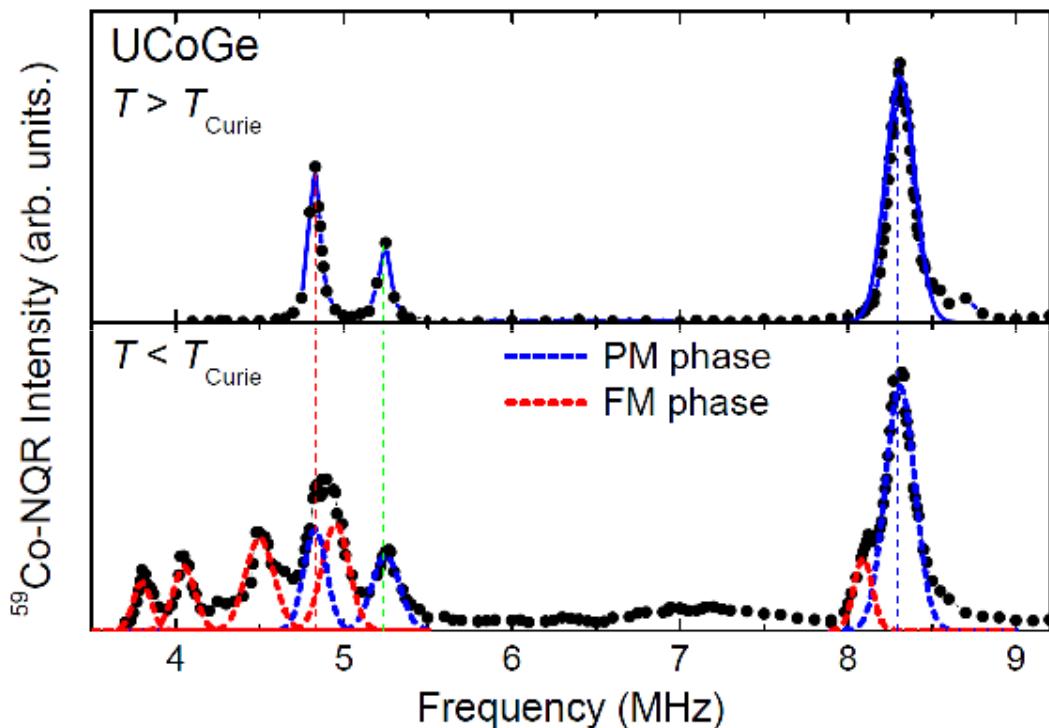


Direction of Magnetic moment



U ordered moments produce the Internal field at the Co site.

^{59}Co -NQR spectrum



FM state ($T < T_{\text{Curie}}$) Two Co signal

- Co signal with internal field

Co signal from FM region

- Co signal without internal field

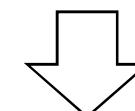
Co signal from PM region



Phase separation

FM region

PM region

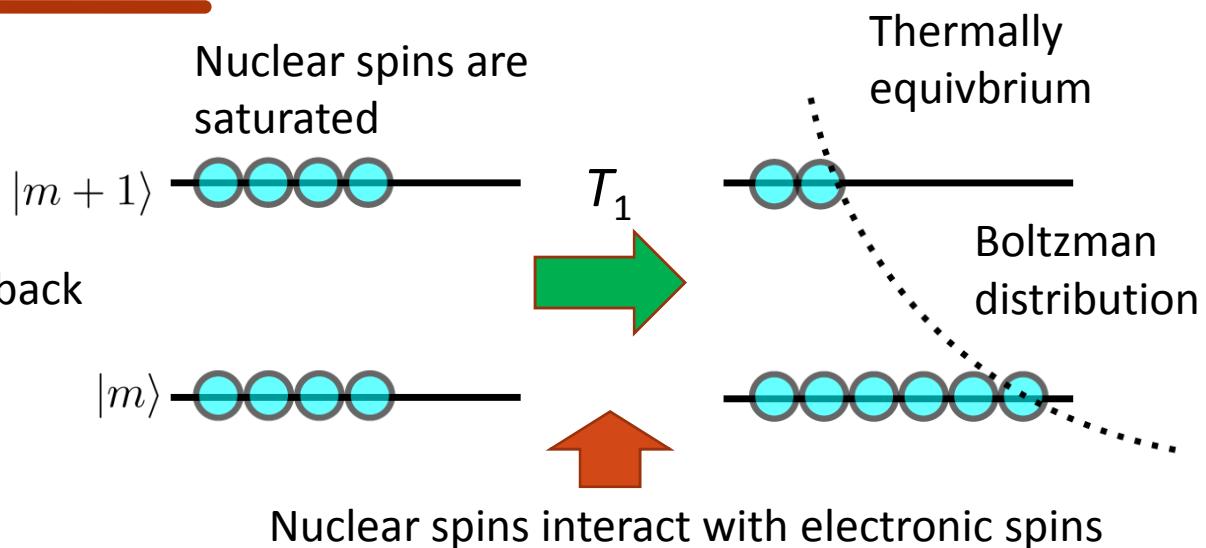


From the measurement of $1/T_1$, we can know which region is responsible for superconductivity.

Spin-lattice relaxation time : T_1

Nuclear Spin-lattice
relaxation Time : T_1

Characteristic time for going back
to the thermal equilibrium
after thermally excited state



$$\frac{1}{T_1} = \frac{2\gamma_n^2 k_B T}{(\gamma_e \hbar)^2} \sum_{\mathbf{q}} A_{\mathbf{q}} A_{-\mathbf{q}} \frac{\chi''(\mathbf{q}, \omega_0)}{\omega_0}$$

$$\cong \frac{4\pi}{\hbar} (\gamma_n \hbar H_{\text{hf}}^s)^2 \frac{N^2(E_F)}{N^2} k_B T$$

In a conventional metal, $1/T_1$ is
proportional to T and $N(E_F)^2$

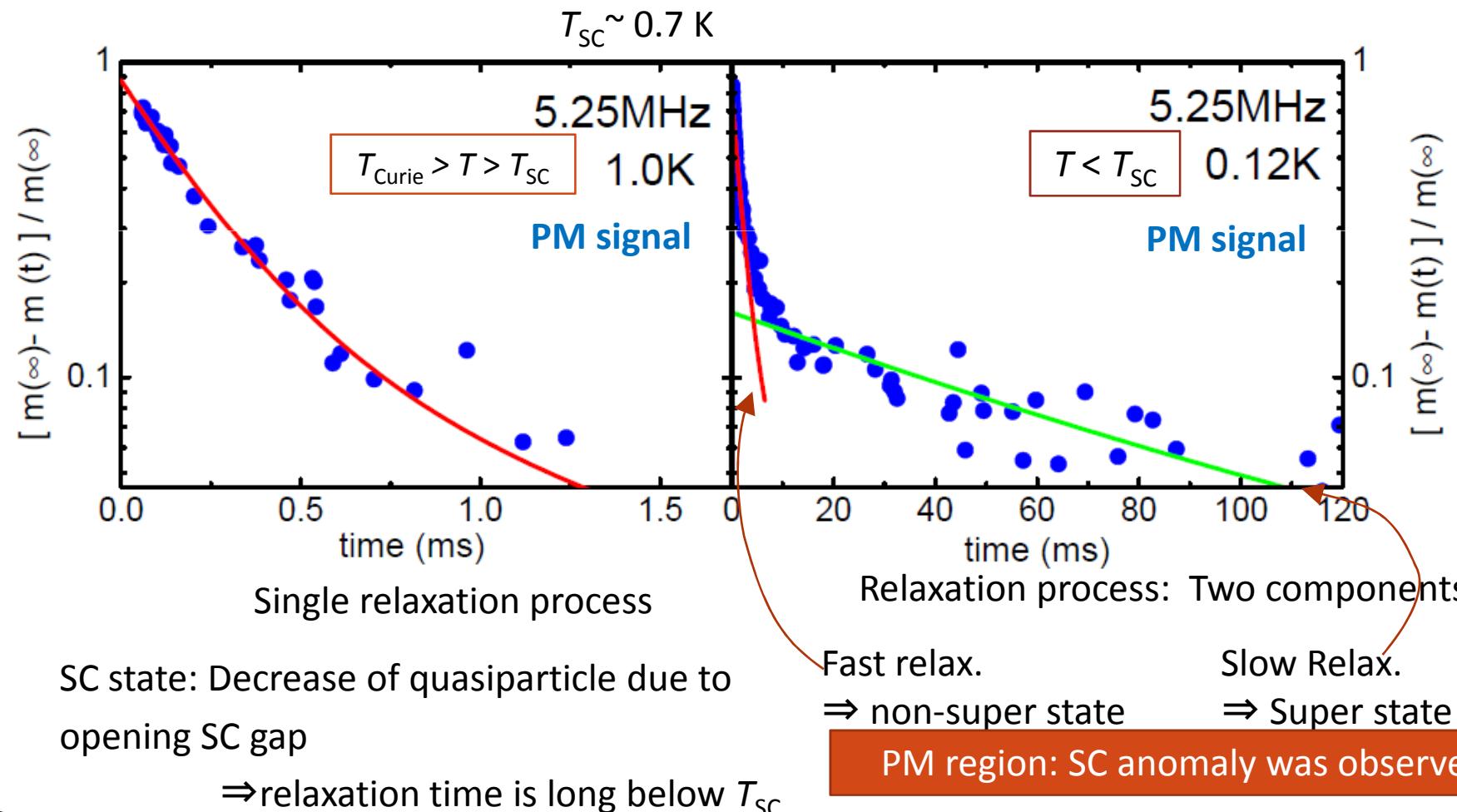
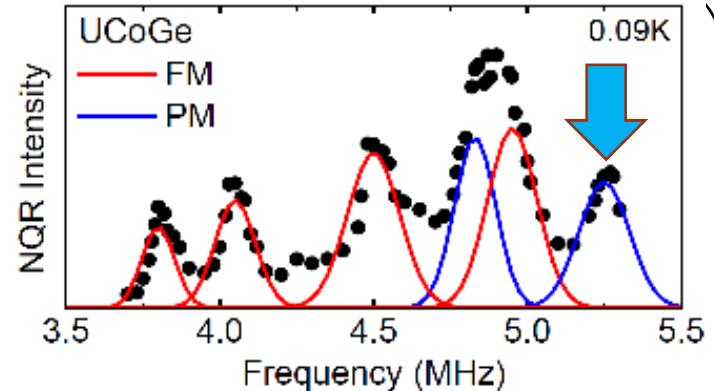
Dynamical information
surrounding the nuclear spins
(e.g. Electrons around EF and
spin dynamics)

Superconducting state
⇒ quasiparticle decreases due to
opening SC gap
⇒ T_1 : longer ($1/T_1$: smaller)

Relaxation Curve @ PM

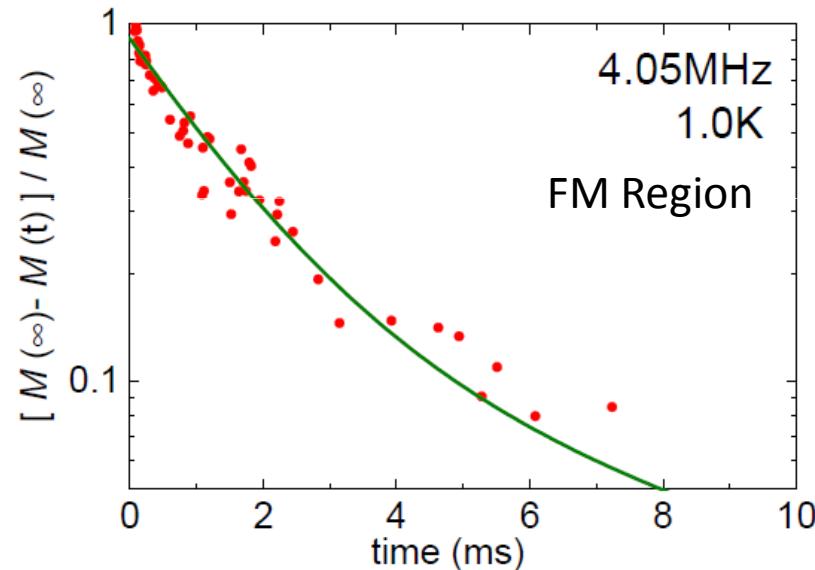
Relaxation behavior : $(m = \pm 5/2 \leftrightarrow \pm 3/2)$

$$\frac{M_\infty - M_z(t)}{M_\infty} \propto \frac{2}{21} \exp\left(-\frac{3t}{T_1}\right) + \frac{25}{154} \exp\left(-\frac{10t}{T_1}\right) + \frac{49}{66} \exp\left(-\frac{21t}{T_1}\right)$$

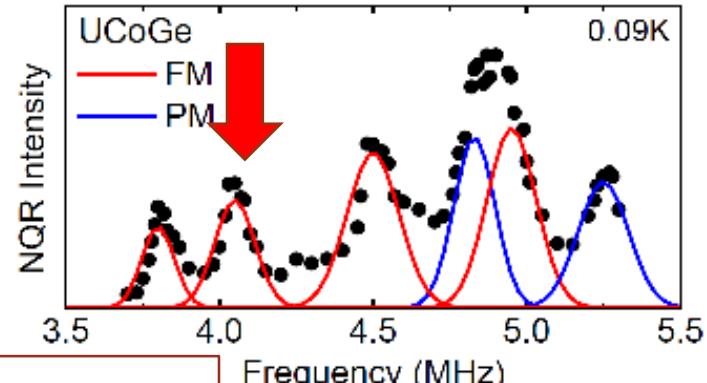


Relaxation Curve @ FM

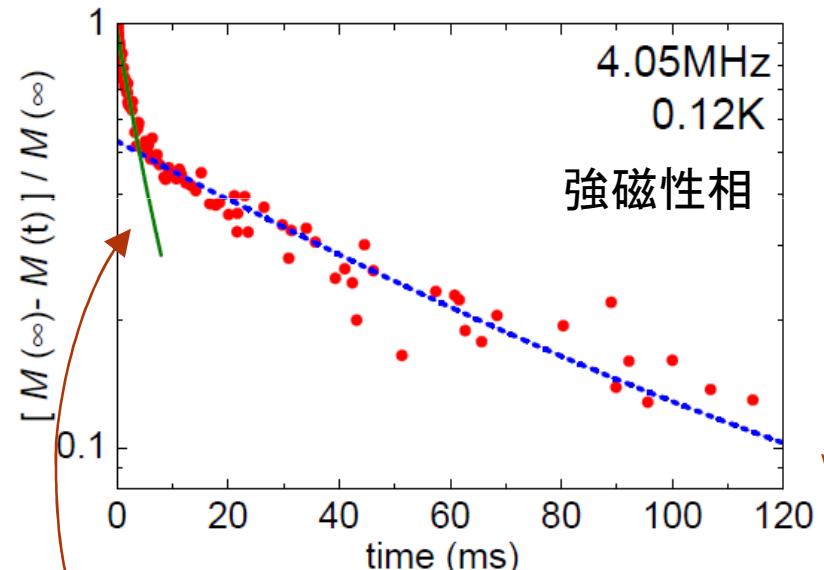
$T > T_{SC}$



Single Relax. Process



$T < T_{SC}$



Relax. Process : Two Comp.
Non-SC state SC state

Microscopic Coexistence of FM and SC ?

Volume fraction of SC

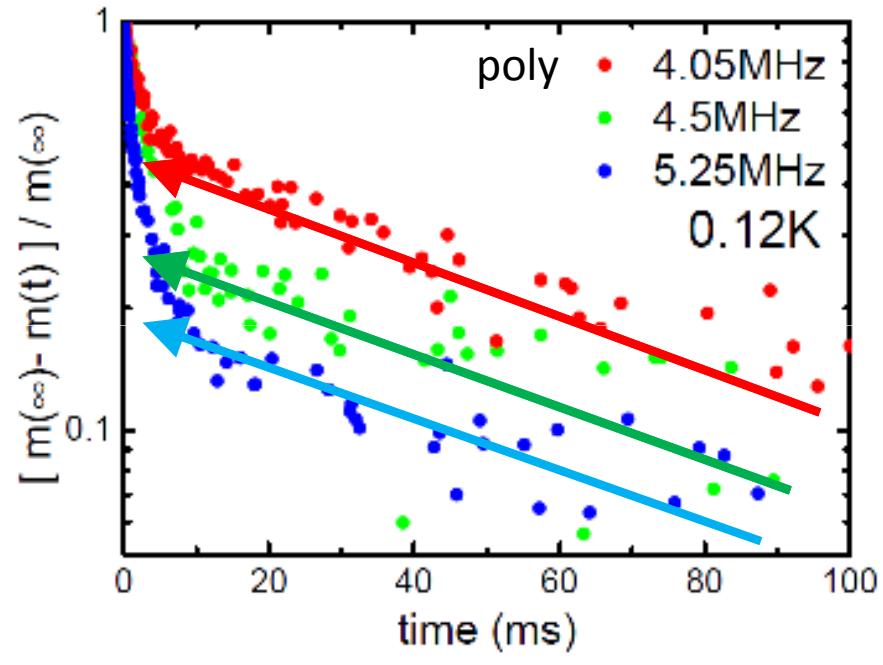
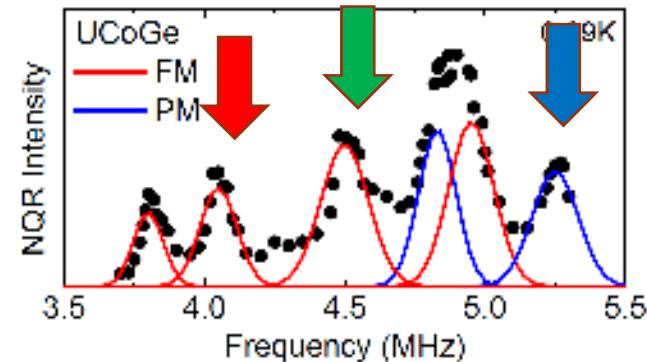
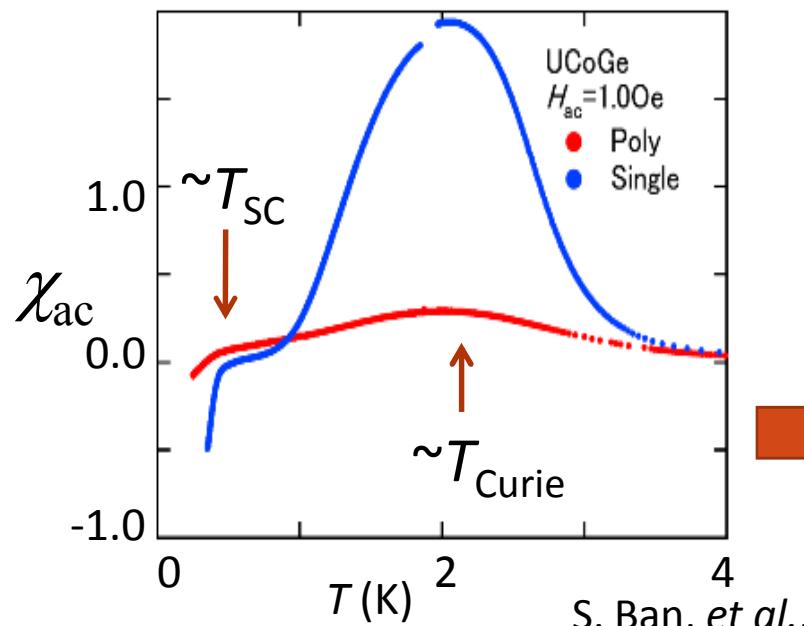
Fraction of SC-state relax. comp. @0.12K

FM ● 4.05MHz ~50%

● 4.5MHz ~30%

PM ● 5.25MHz ~20%

Fraction of SC state SC comp. at the FM signal is larger than that at the PM signal



Anomaly at T_{curie} is obvious
⇒ Anomaly at T_{SC} is sharp

Intimate relation between FM and SC states