

Anomalous Proximity Effect in Triplet Superconductor Junctions

Yukio Tanaka

**Department of Applied Physics
Nagoya University**

2005 12/15 「スーパークリーン物質で実現する新しい量子相の物理」

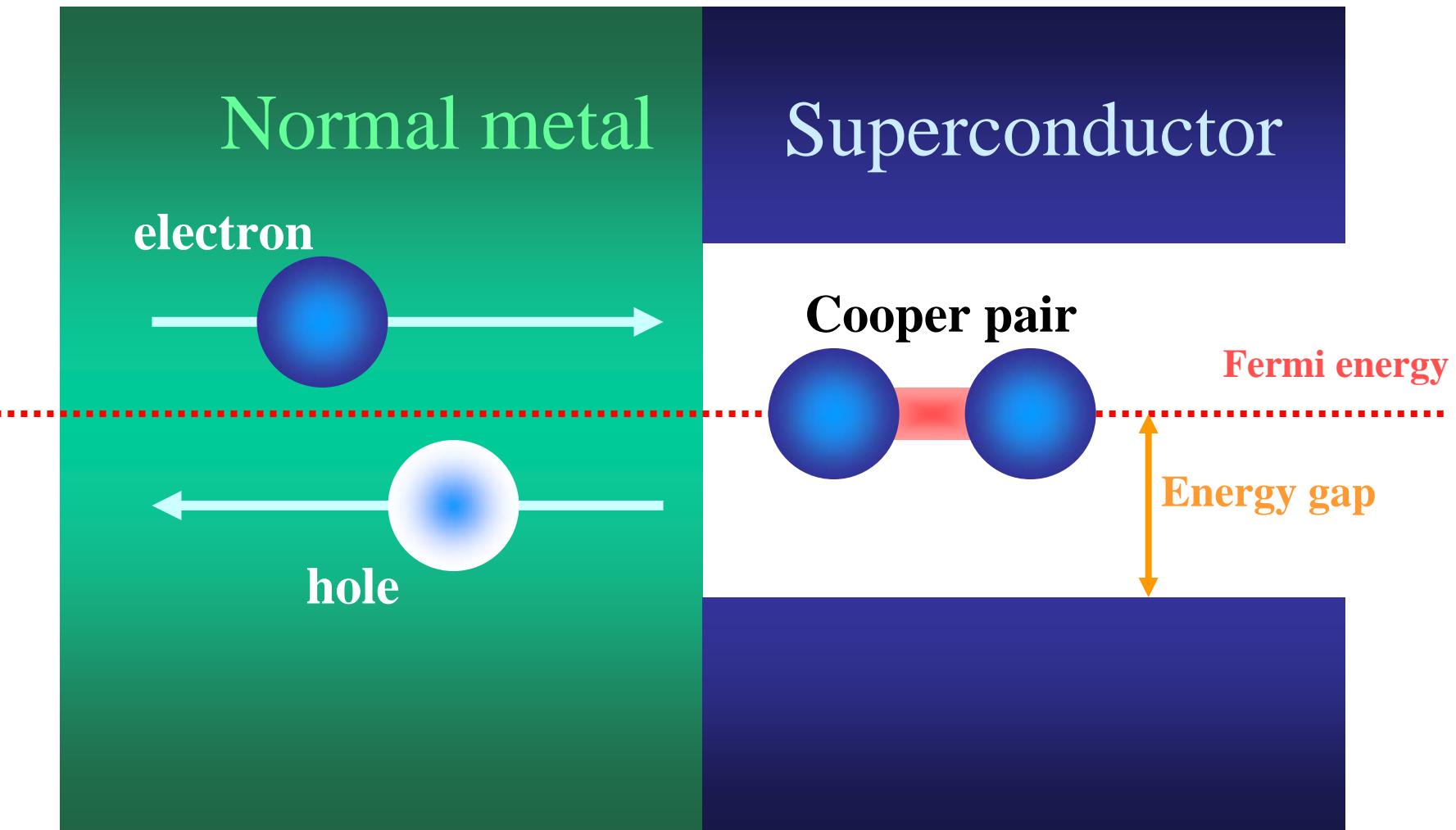
Main Collaborators

- **S. Kashiwaya**
**National Institute of Advanced
Industrial Science and Technology**
- **A. Golubov** Twente University
- **Y. V. Nazarov** Delft University
- **Y. Asano** Hokkaido University
- **T. Yokoyama** Nagoya University

Contents of this talk

- (1) Mid gap Andreev resonant state
- (2) Proximity effect in singlet superconductor junctions [PRB 69 144519 (2004), PRL 90 167003(2003)]
- (3) Proximity effect in triplet superconductor junctions [PRB 70, 012507 (2004), PRB71 024506(2005)]
- (4) Future Problems

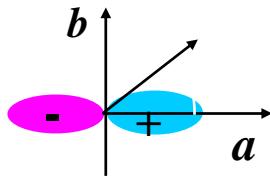
Andreev reflection



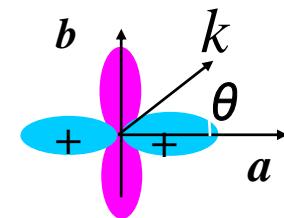
1964 Andreev

Tunneling spectroscopy of Unconventional superconductors

p_x - wave



$d_{x^2-y^2}$ -wave



$$\Delta(\theta) = \Delta_0 \cos \theta$$

$$\Delta(\theta) = \Delta_0 \cos 2\theta$$

Triplet superconductor

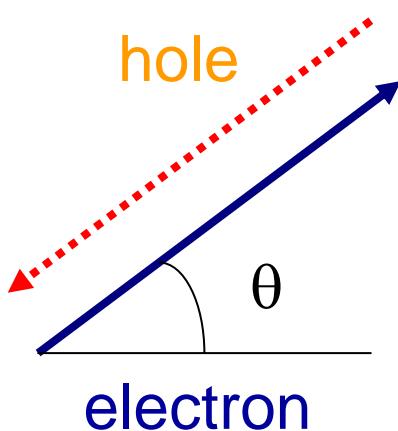
Singlet unconventional superconductor

Quasiparticles feel **different sign of the pair potential** depending on their motions.

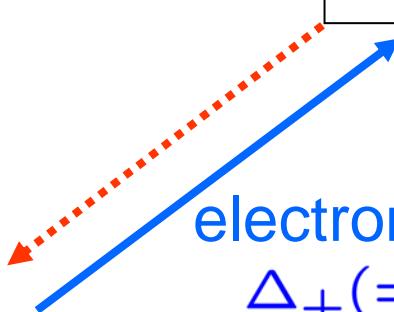
What is expected in unconventional superconductor junctions?

Mid gap Andreev resonant state (MARS)

Normal metal



$$\Delta_+ \Delta_- < 0$$

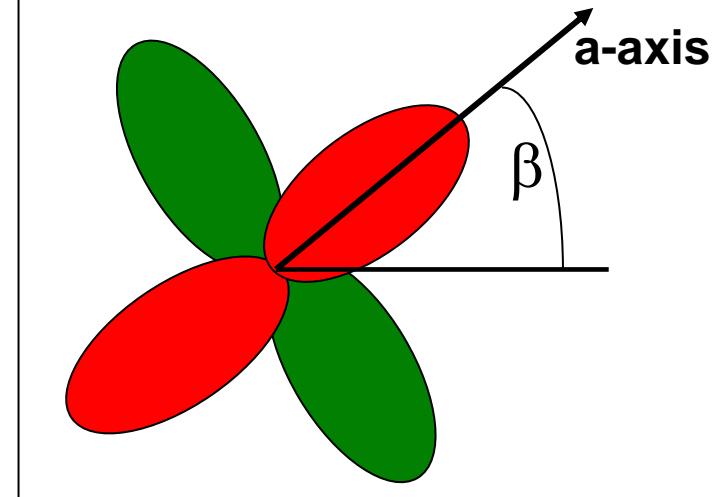


$$\Delta_+ (= \Delta_0 \cos[2(\theta - \beta)])$$

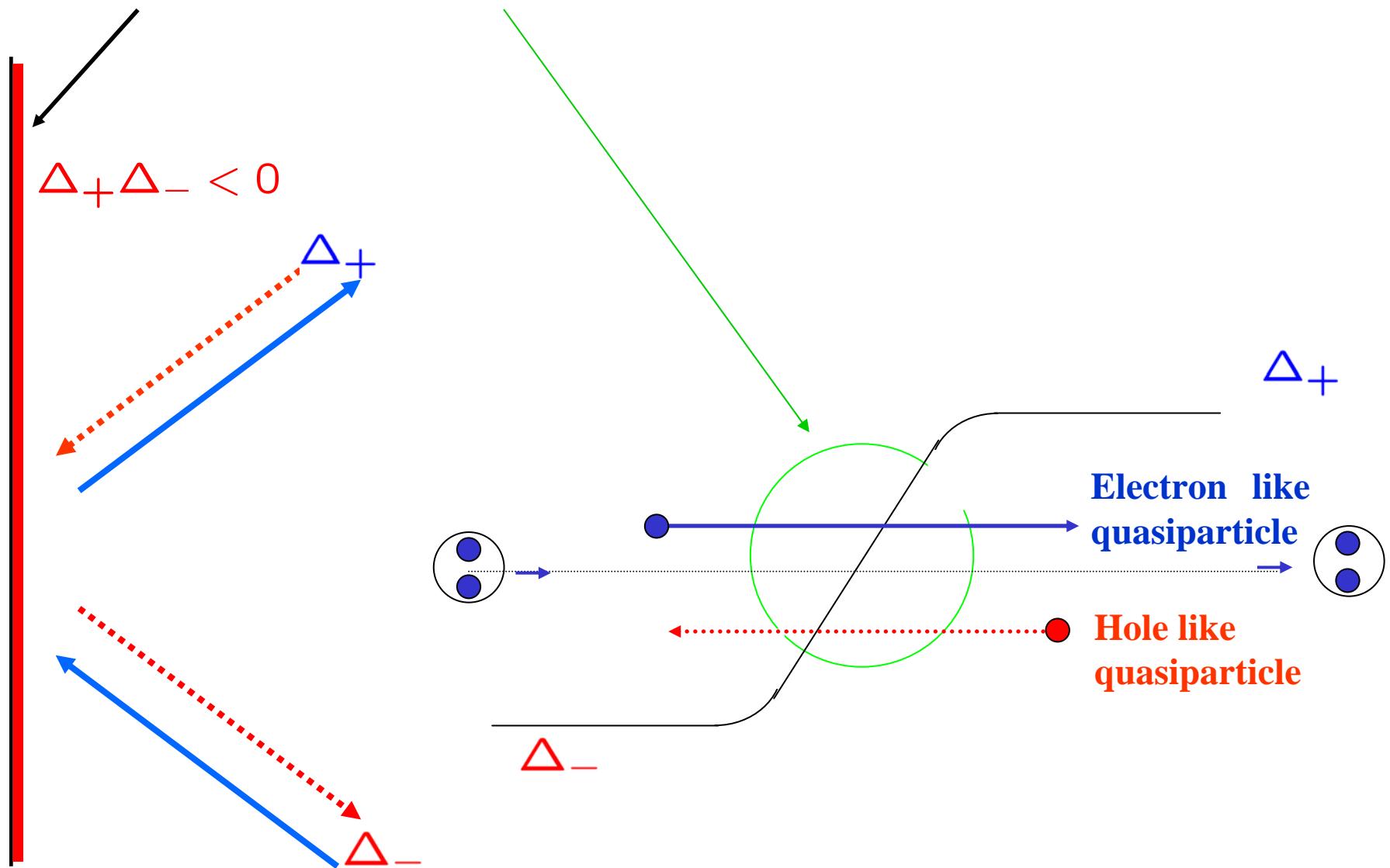
hole like quasiparticle

$$\Delta_- (= \Delta_0 \cos[2(\theta + \beta)])$$

Unconventional
superconductor

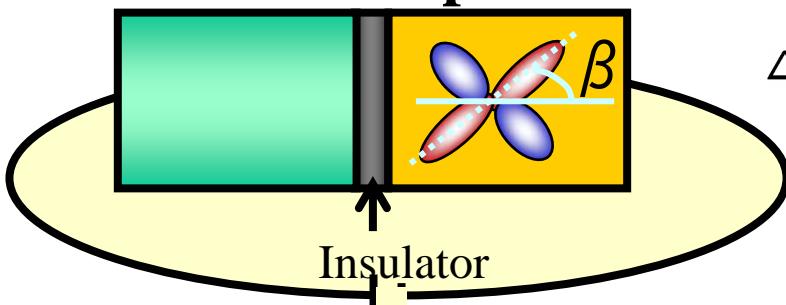
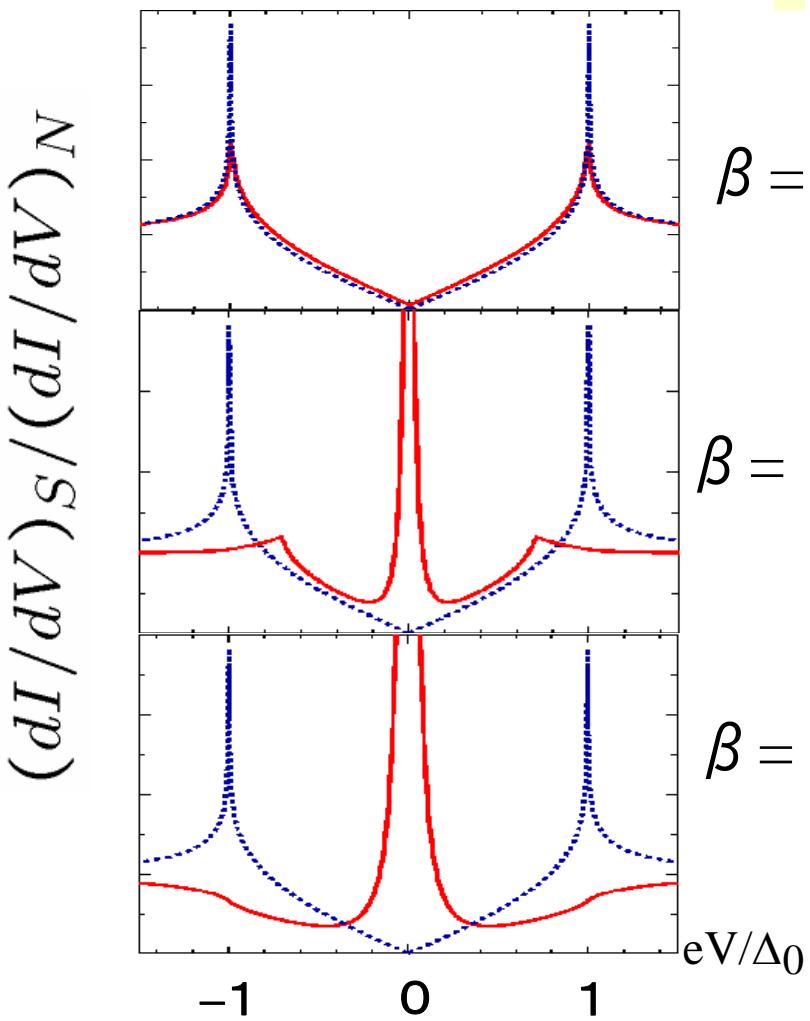


Mid gap Andreev resonant state (MARS)



Tunneling conductance in d-wave superconductor junction (ballistic)

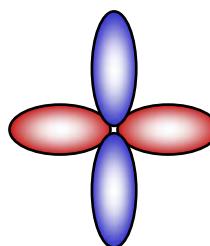
Blue dotted line:
Bulk d-wave DOS



$$\Delta_{\pm} = \Delta_0 \cos[2(\theta \mp \beta)]$$

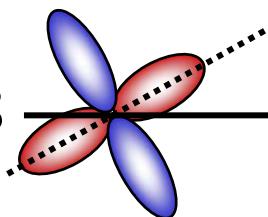
Bruder (1990)
Blonder Tinkham
Klapwijk (1982)

$$\beta = 0$$



Tanaka Kashiwaya (1995)
Phys Rev Lett 74 3451 (1995)

$$\beta = \pi/8$$

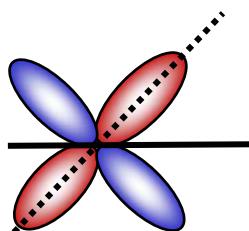


ZBCP

Zero bias conductance peak

Mid gap Andreev
resonant state

$$\beta = \pi/4$$



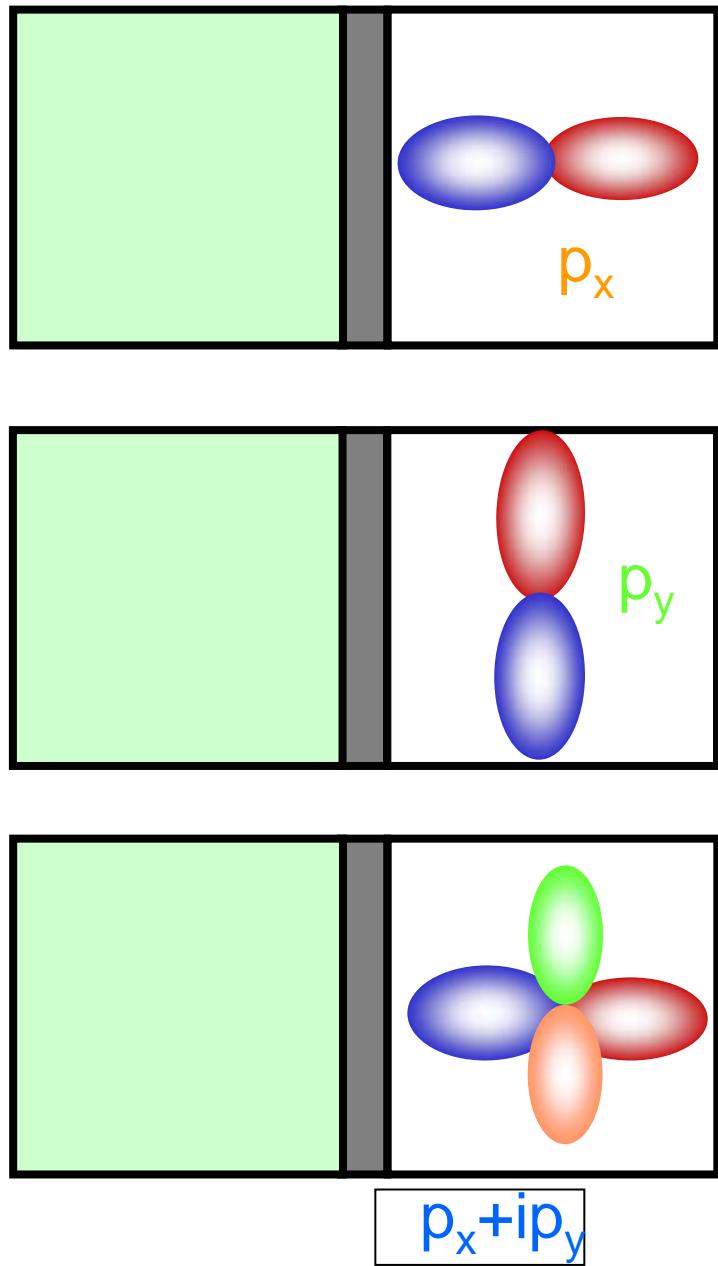
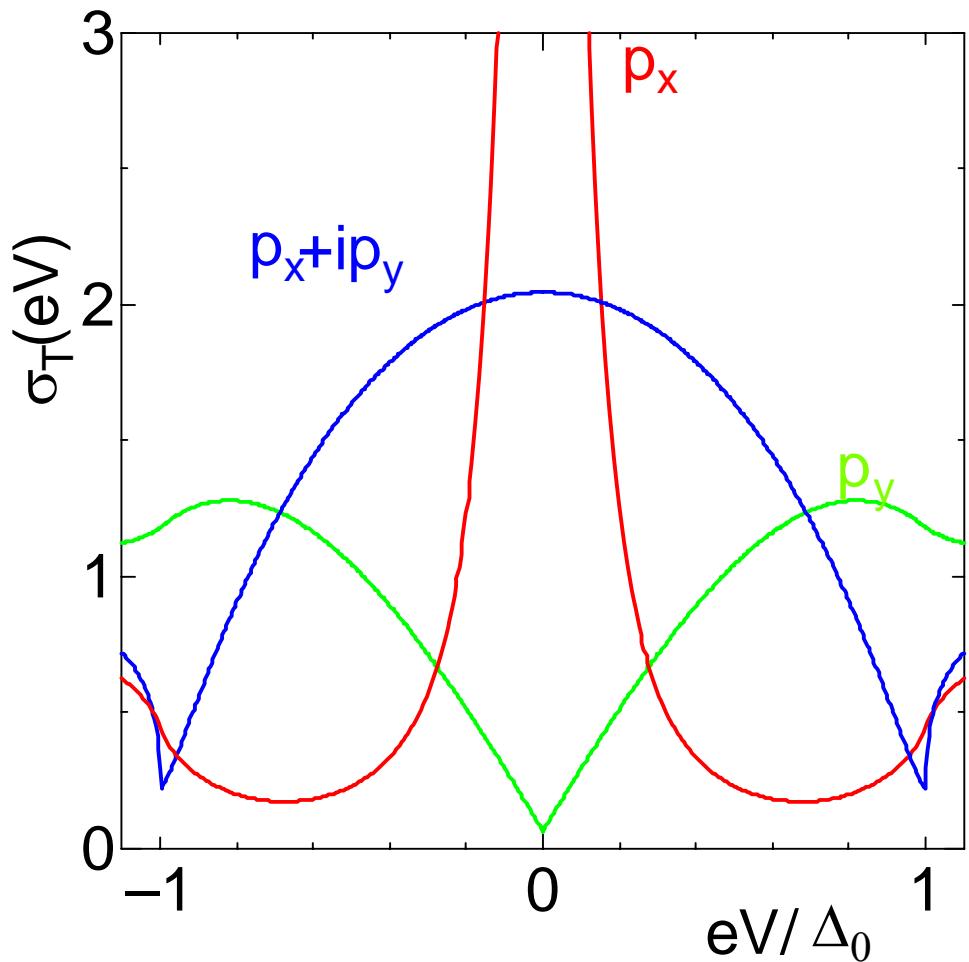
Surface bound state
Hu (1994)
Buchholtz (1981)
Hara Nagai(1986)
Matsumoto Shiba(1995)

Normal metal /triplet junctions

Phys. Rev. B. 56, 7847 (1997)

J. Phys. Soc. Jpn. 71, 2102 (2002)

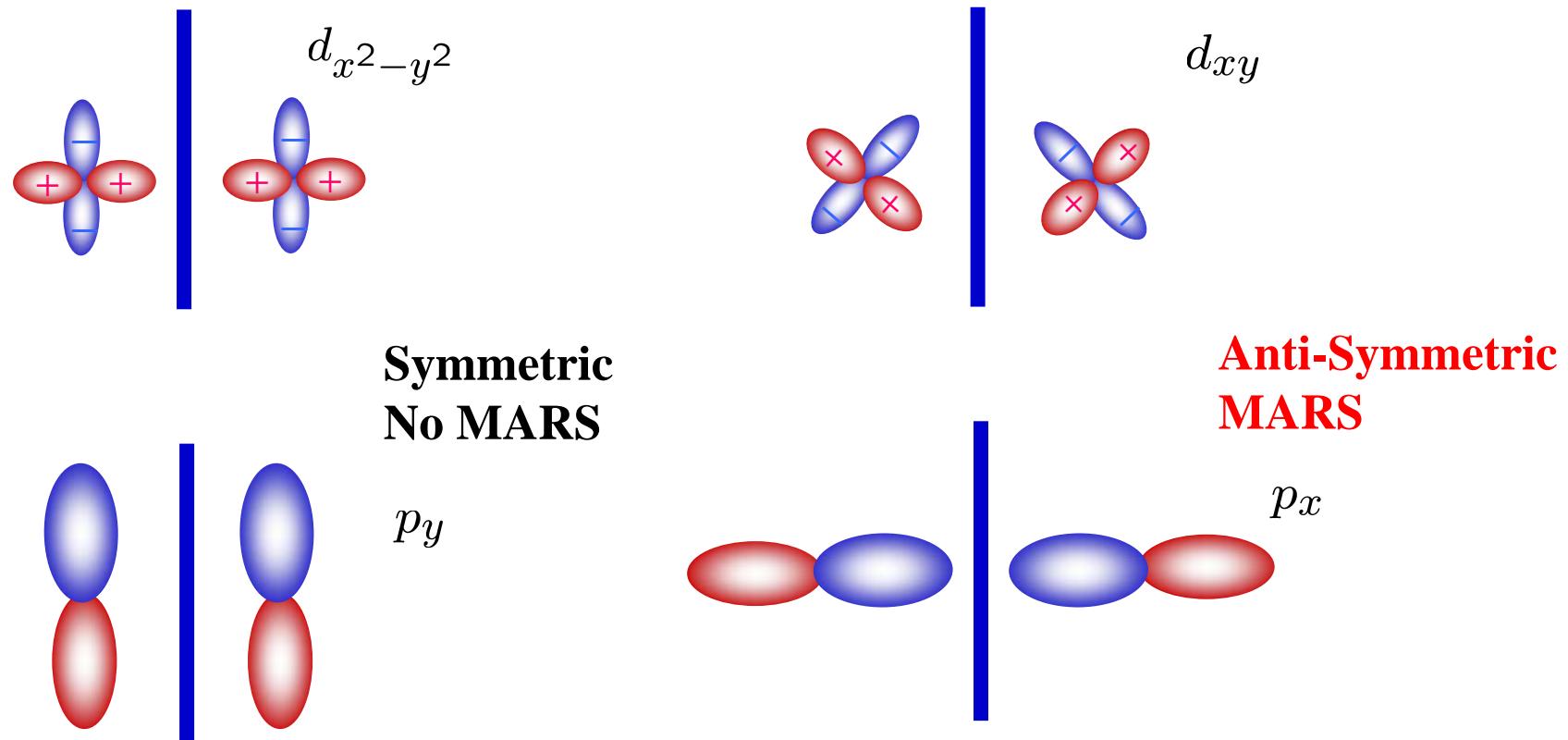
J. Phys. Soc. Jpn. 67, 3224 (1998)



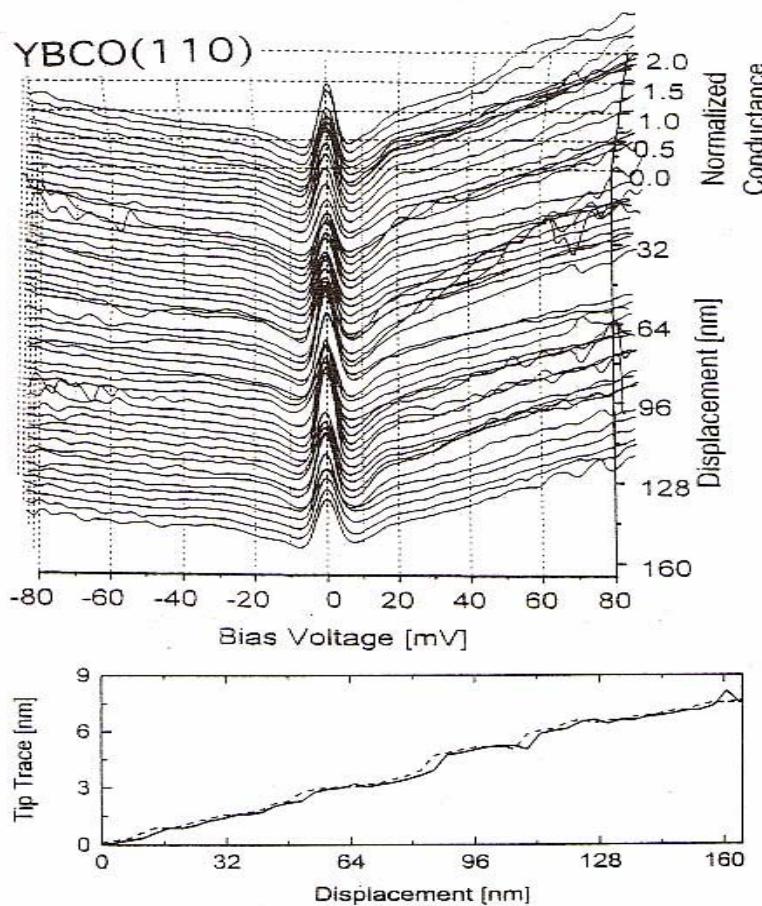
$$\Delta_{\uparrow\downarrow}(\theta) = \Delta_0 \exp(i\theta)$$

Condition of the formation of mid gap Andreev resonant state(MARS)

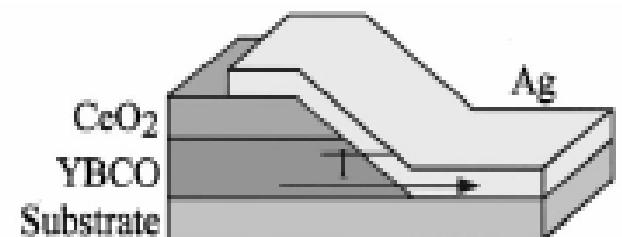
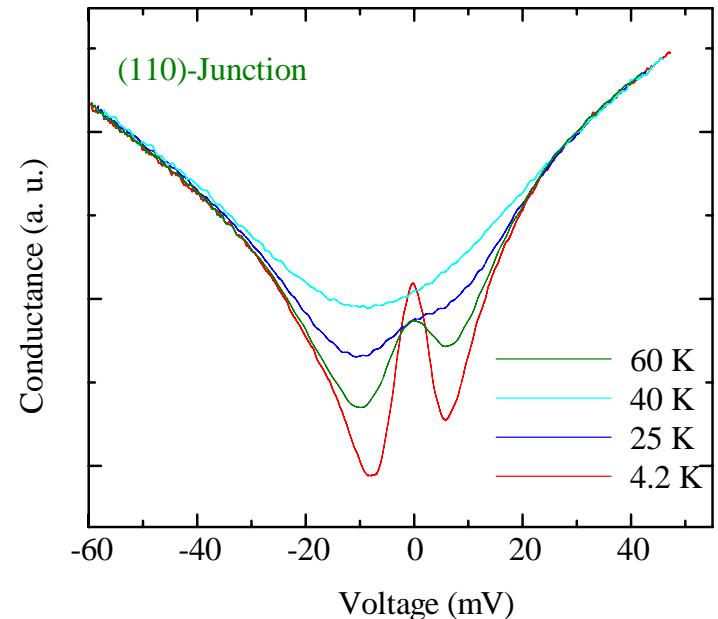
Inversion at the plane parallel to the interface



Zero bias conductance peak (by MARS) observed in cuprate



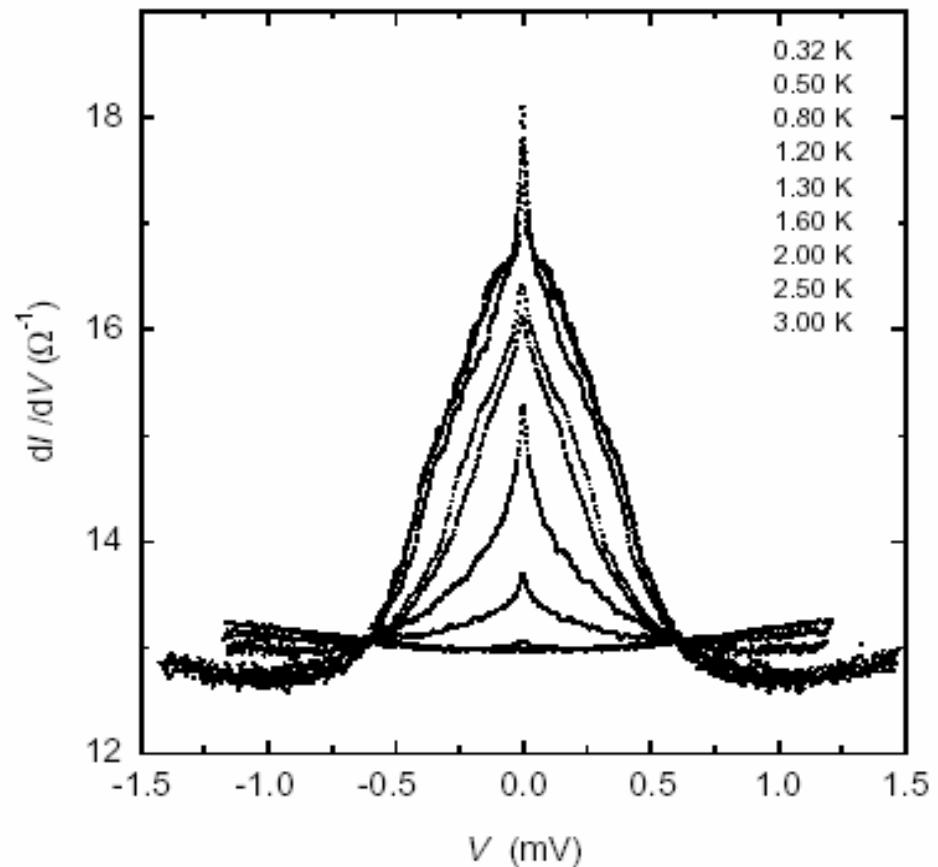
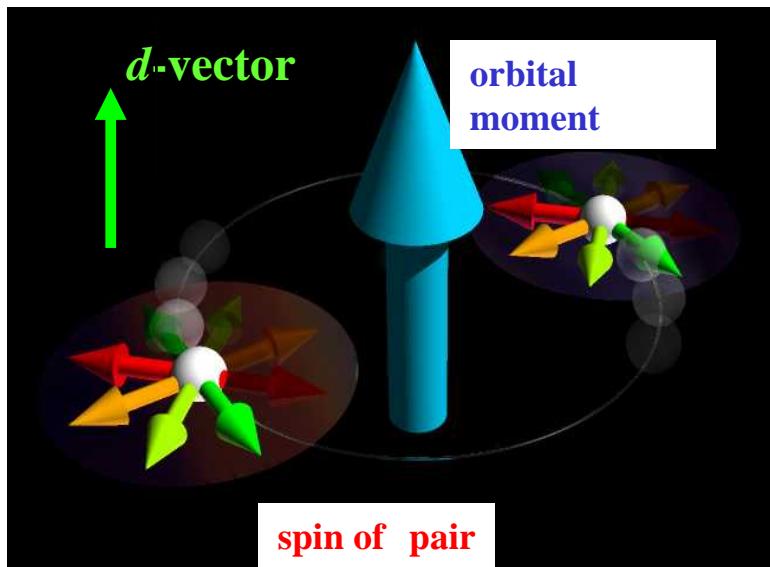
S. Kashiwaya and Y. Tanaka
Rep. Prog. Phys. (2000)



Iguchi Wang et al. –
Phys. Rev. B60, 4272 (1999)

MARS observed in triplet superconductor Sr_2RuO_4

Y. Maeno, G. Bednorz et al.
Nature 372 532 (1994)
(p-wave chiral)



Mao, Nelson, Jin, Liu and Maeno Phys. Rev. Lett. 87, 037003 (2001)
Kawamura, Yaguchi, Kikugawa Maeno Takayanagi
J. Phys. Soc. Jpn. 74 531 (2005)

Superconducting Materials where MARS is observed

$\text{YBa}_2\text{CuO}_{7-\delta}$ (Geerk, Kashiwaya, Iguchi, Greene, Yeh, Wei...)

$\text{Bi}_2\text{Sr}_2\text{CaCu}_2\text{O}_y$ (Ng, Suzuki, Greene....)

$\text{La}_{2-x}\text{Sr}_x\text{CuO}_4$ (Iguchi)

$\text{La}_{2-x}\text{Ce}_x\text{CuO}_4$ (Cheska)

$\text{Pr}_{2-x}\text{Ce}_x\text{CuO}_4$ (R.L.Greene)

Sr_2RuO_4 (Mao, Meno, Kawamura,Lube)

$\kappa-(\text{BEDT-TTF})_2X$, $X=\text{Cu}[\text{N}(\text{CN})_2]\text{Br}$ (Ichimura)

UBe_{13} (Ott)

CeCoIn_5 (Wei)

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- (4) Future Problems

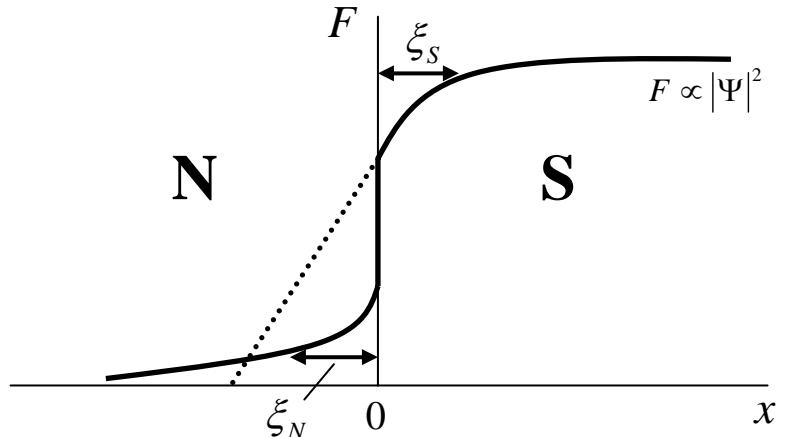
Proximity effect and total resistance of the junction

Pair amplitude

$$F = \langle \Psi_{\uparrow} \Psi_{\downarrow} \rangle$$

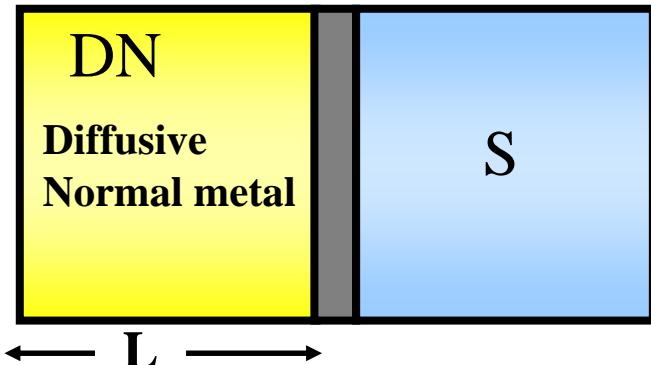
Normal metal

$$F \propto \exp\left(-\frac{|x|}{\xi_N}\right)$$



**Proximity effect has an crucial influence on the charge transport
in diffusive normal metal / s-wave conventional superconductor junction
(DN / S junction)**

insulator



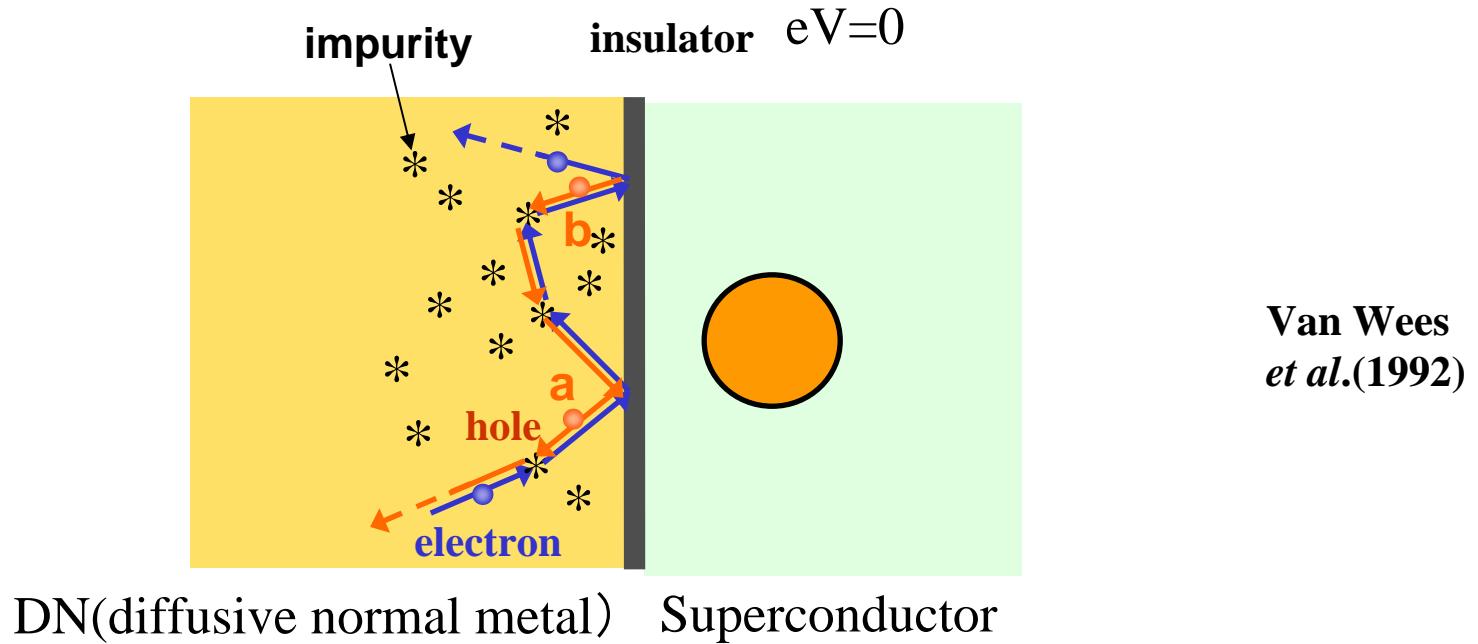
ξ : localization length

l : mean free path

$$l \ll L \ll \xi$$

Retro reflectivity of the Andreev reflection in DN

Beenakker: Rev. Mod. Phys. (97)

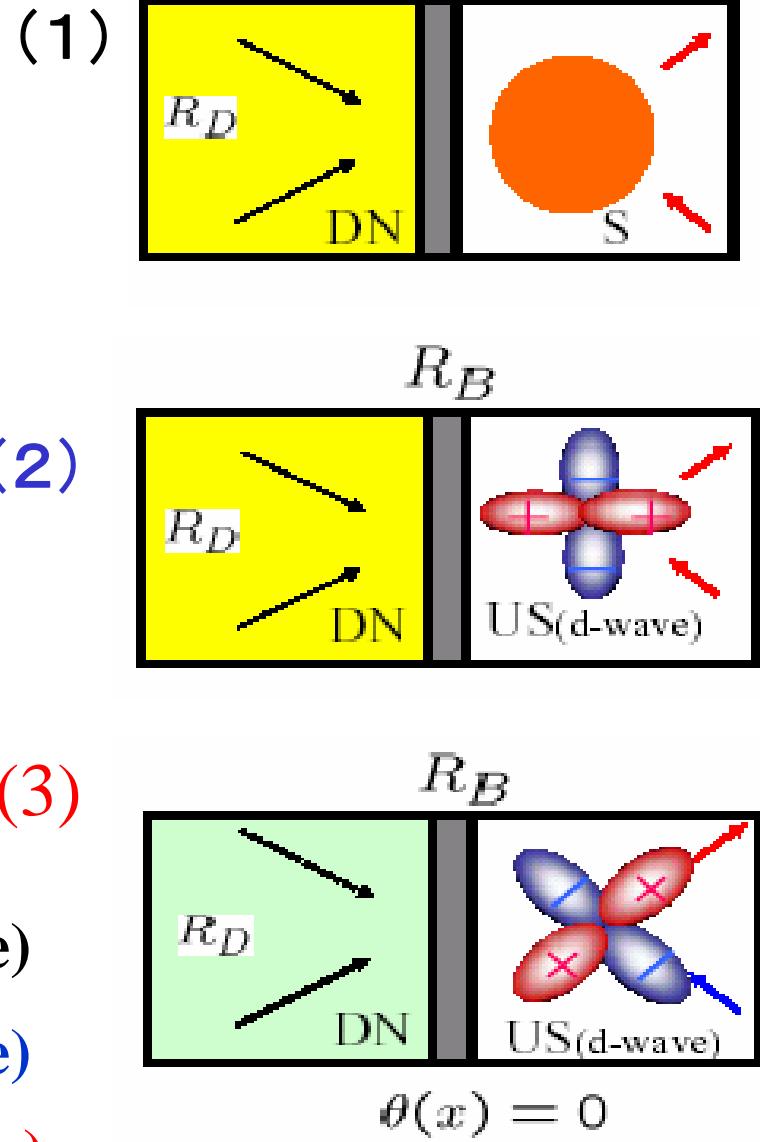
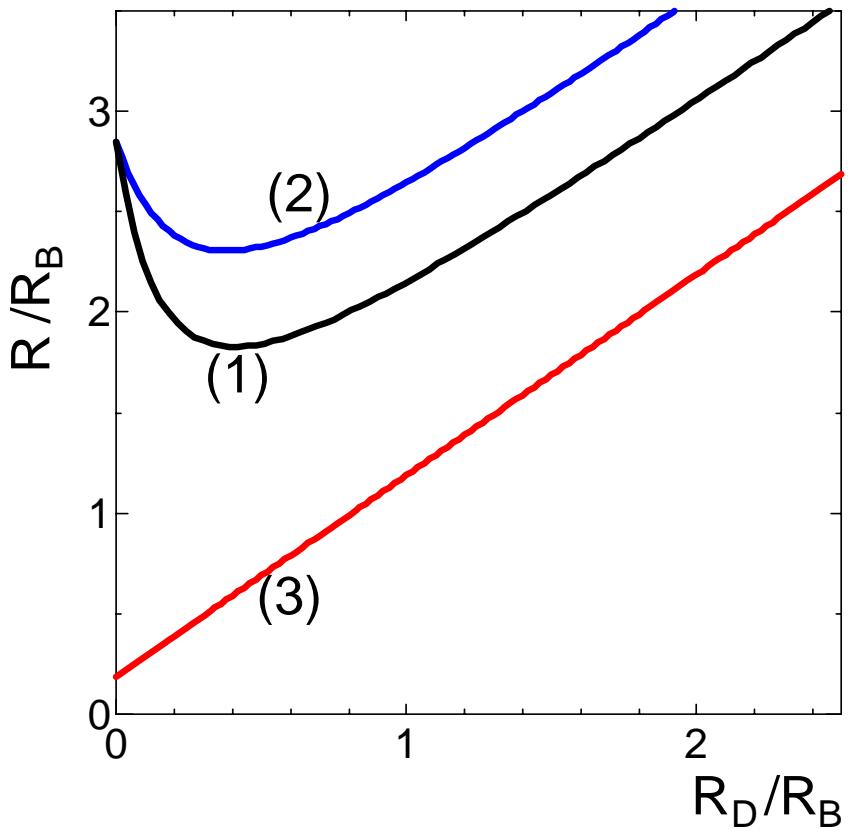


Van Wees
et al.(1992)

Proximity effect is enhanced by the Andreev reflection

It is effective to use quasiclassical Green's function to describe charge transport.

Zero voltage total resistance

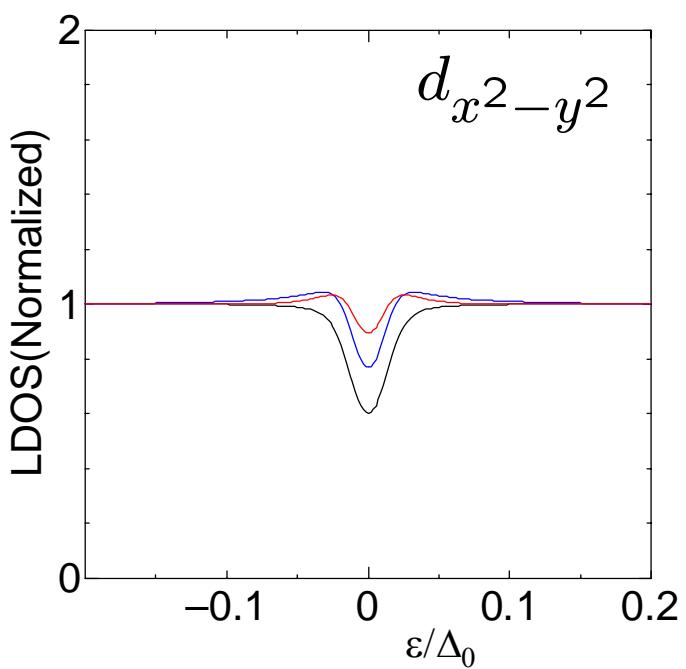
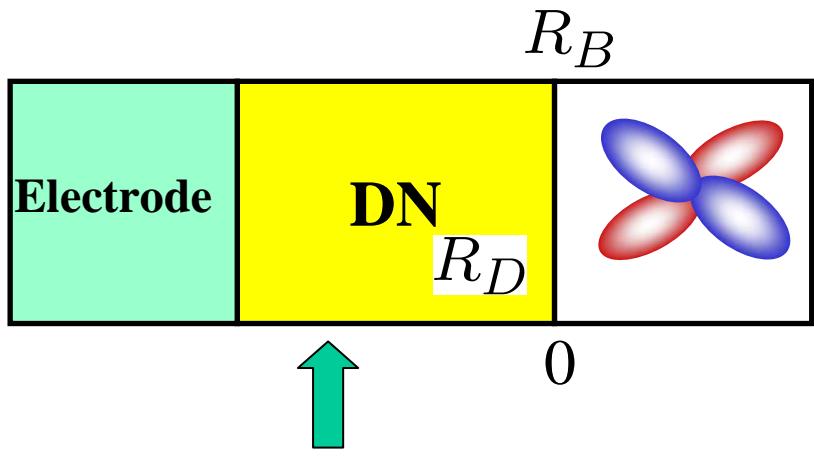
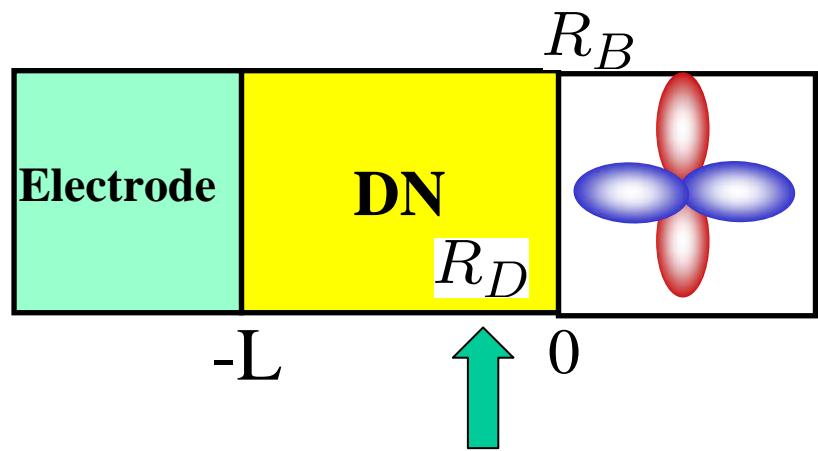


- (1) Proximity no MARS (reentrance)
- (2) Proximity no MARS (reentrance)
- (3) No Proximity MARS (Ohm's rule)

R_0 ; Sharvin resistance

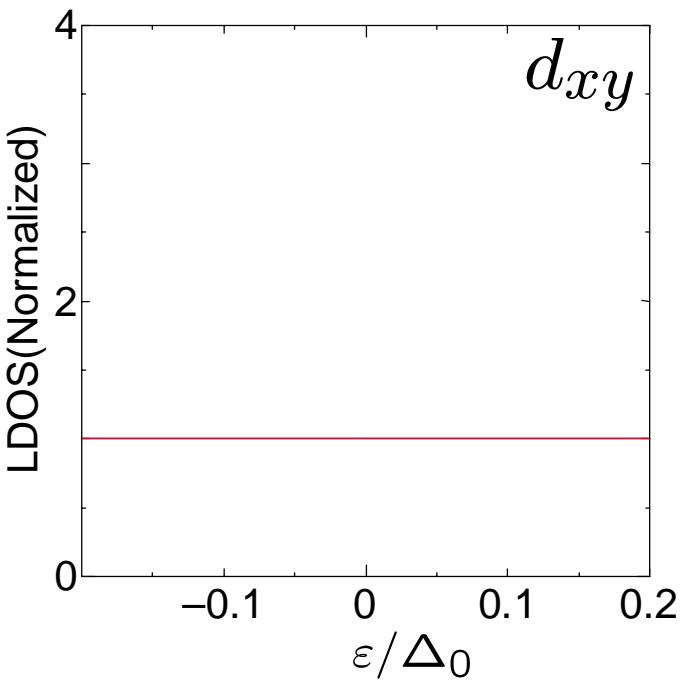
$$R_B = 2R_0 / \int_{-\pi/2}^{\pi/2} T(\phi) \cos \phi d\phi$$

Local density of states in DN (d-wave)



Energy gap due to proximity effect

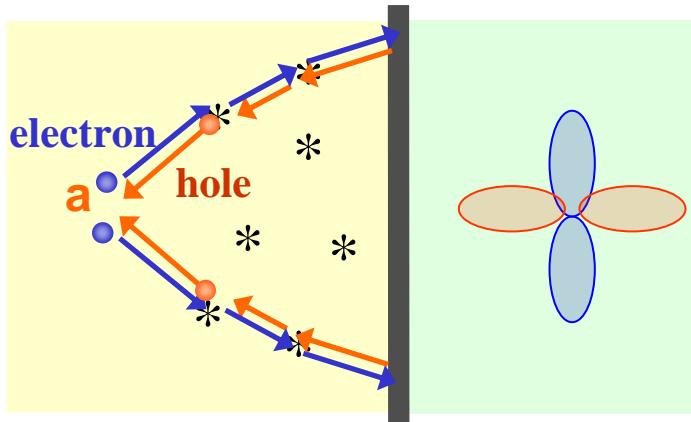
$Z=10$
 $L/\xi = 18$
 $E_{th}/\Delta_0 = 0.01$
 $\xi = \sqrt{D/(2\pi T_C)}$
x=0
x= -L/4
x= -L/2



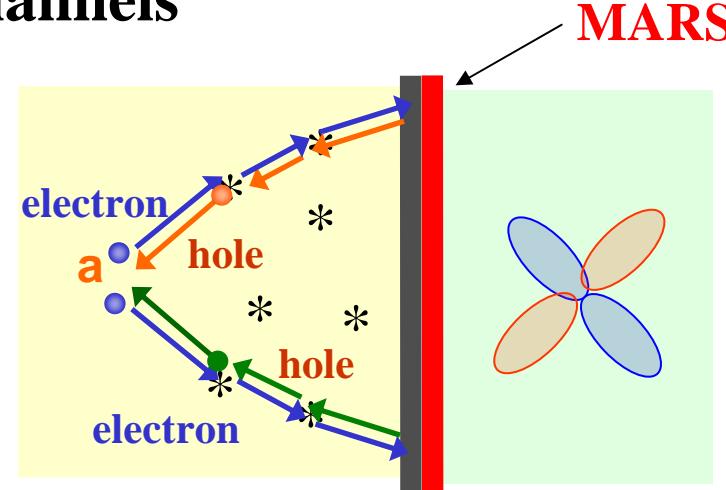
No proximity effect

Condition for the proximity effect

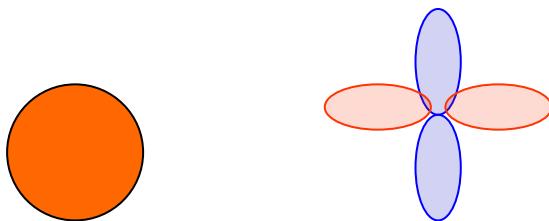
Averaging from many channels



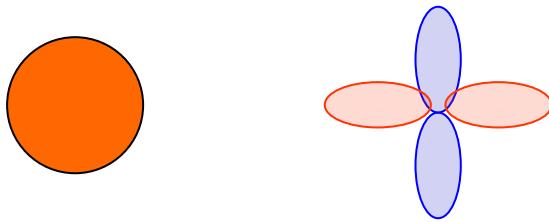
$d_{x^2-y^2}$ Finite proximity



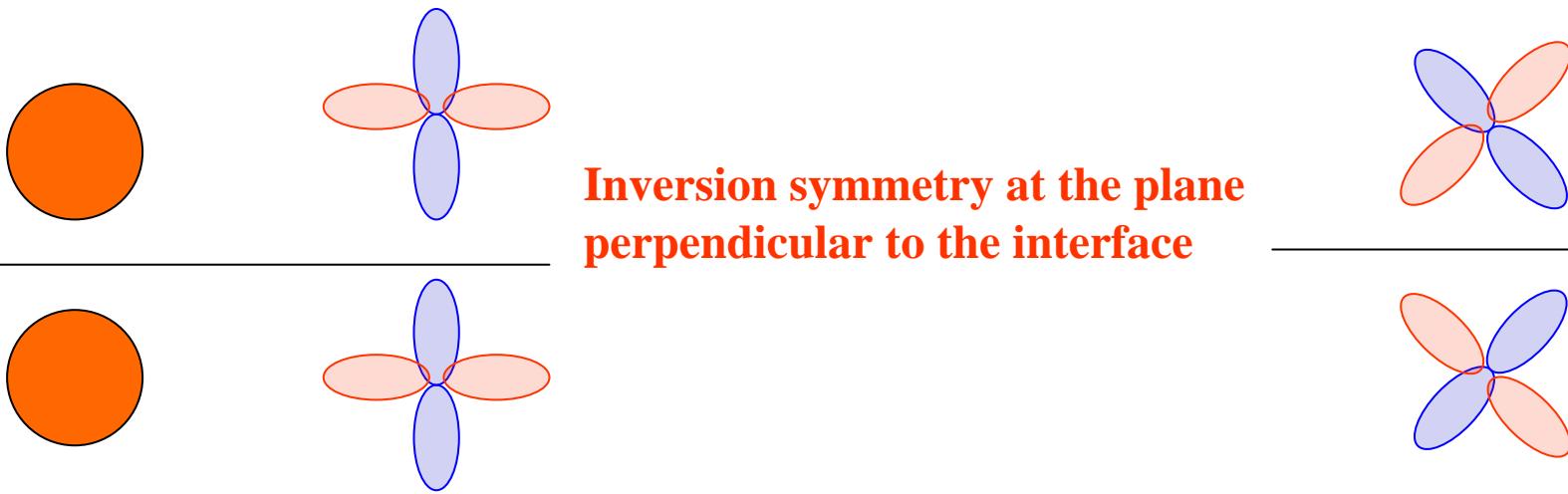
proximity effect is absent



Inversion symmetry at the plane
perpendicular to the interface

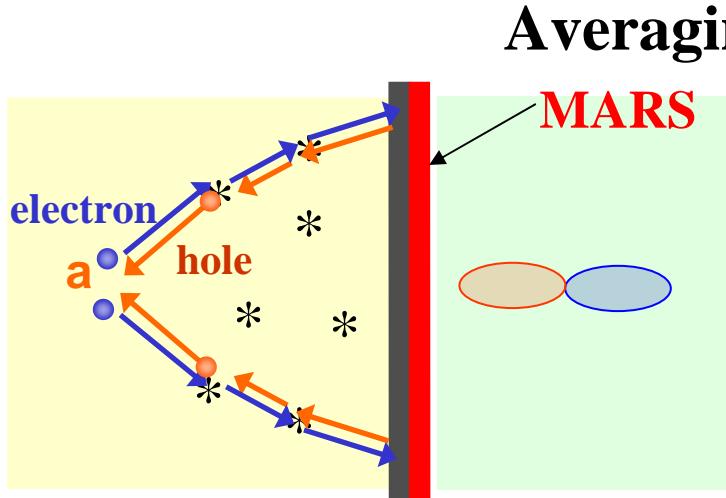
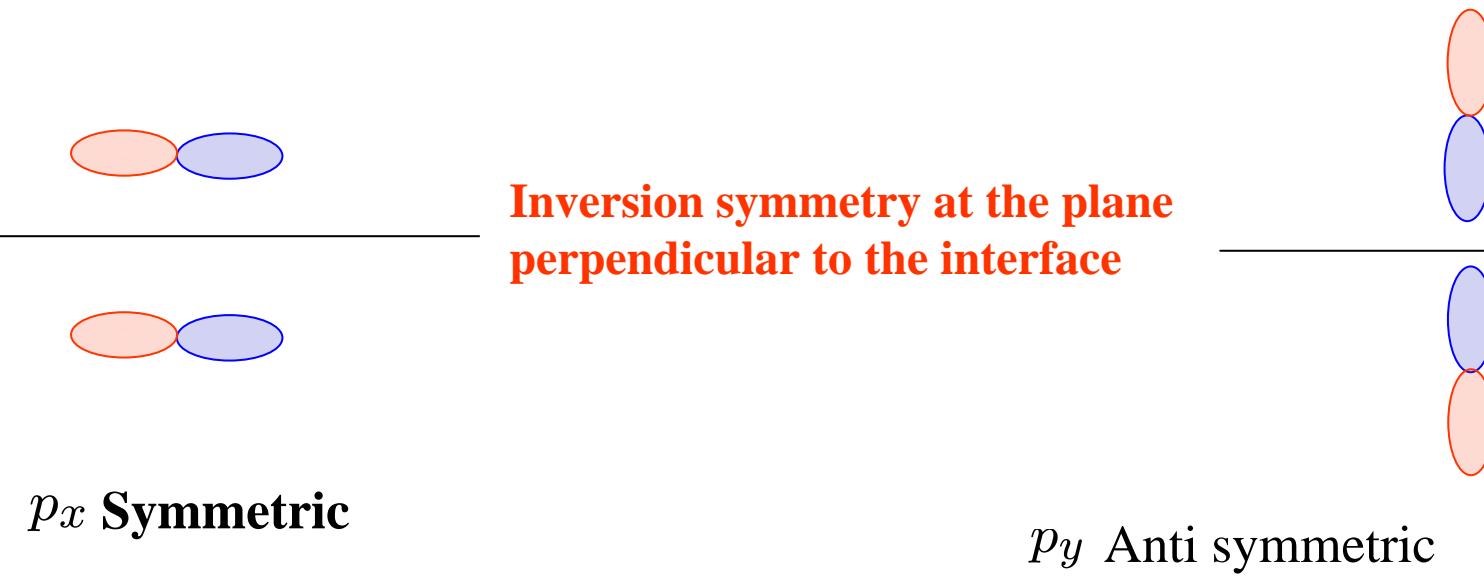


Symmetric

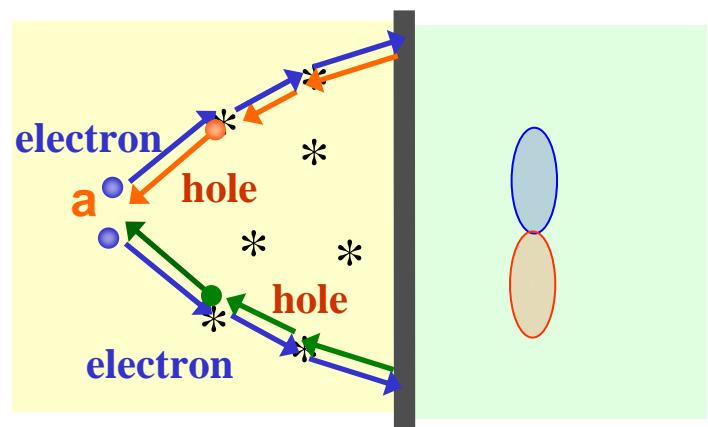


Anti symmetric

Condition for the proximity effect



Finite proximity



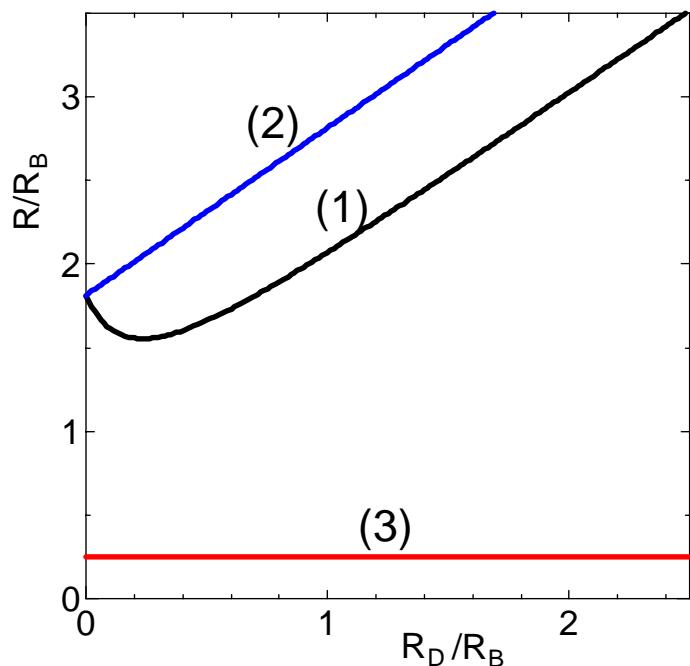
proximity effect is absent

Averaging from many channels

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Total resistance (eV=0)



R_D Resistance in DN

R_B Resistance at the interface

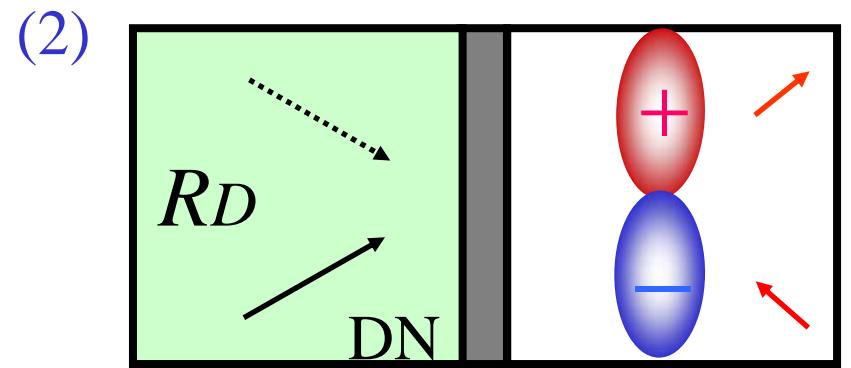
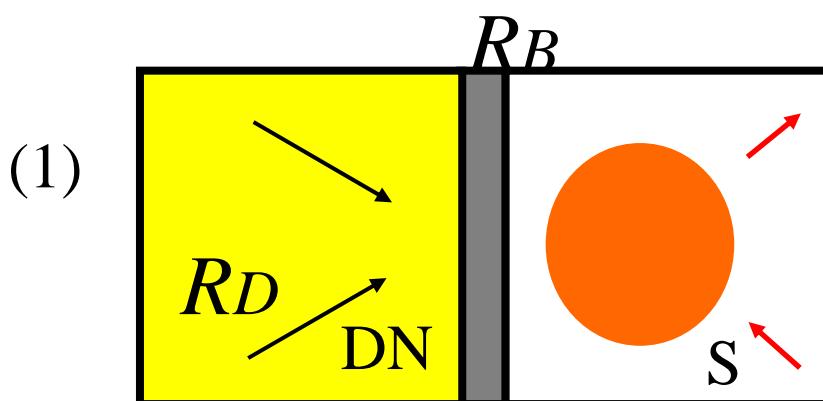
$$(2) \text{ p}_y\text{-wave} \quad R = R_D + R_{R_D=0}$$

(3) $\text{p}_x\text{-wave}$

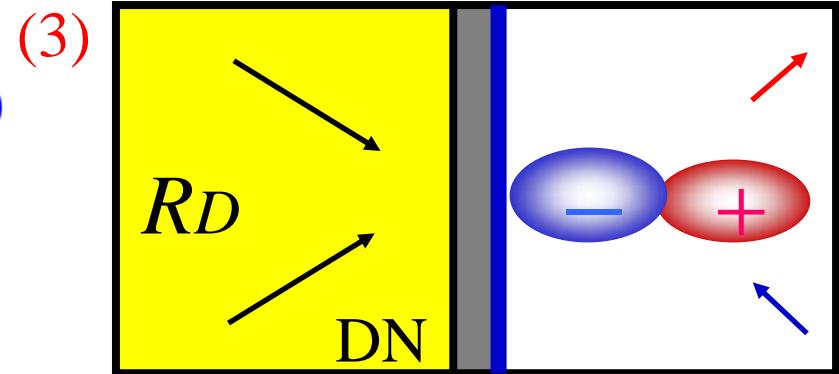
R is independent of R_D !!

$$R = R_0/2$$

$$R_B = 2R_0 / \int_{-\pi/2}^{\pi/2} T(\phi) \cos \phi d\phi \quad R_0; \text{ Sharvin resistance}$$

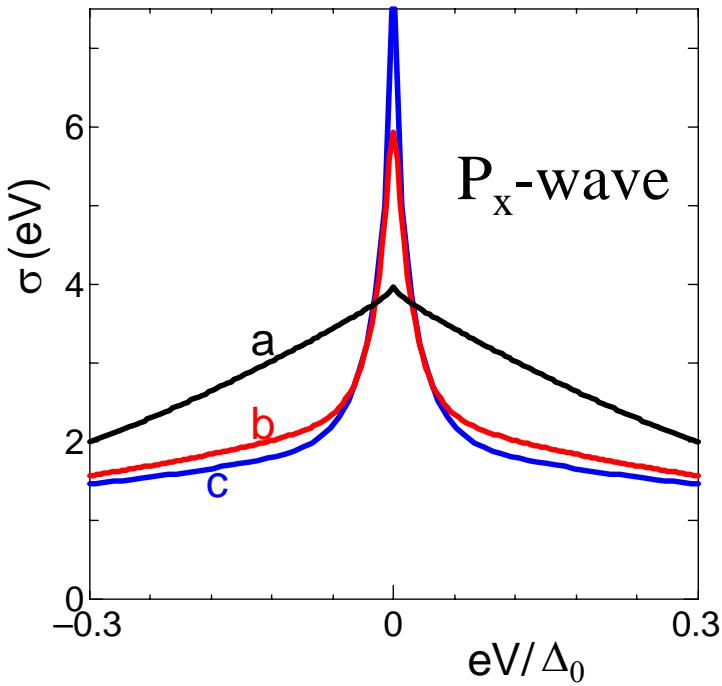


(no proximity & no MARS)

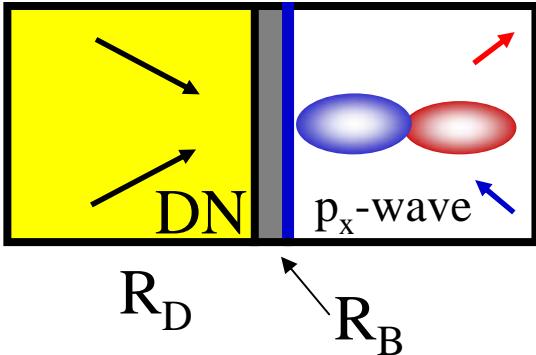
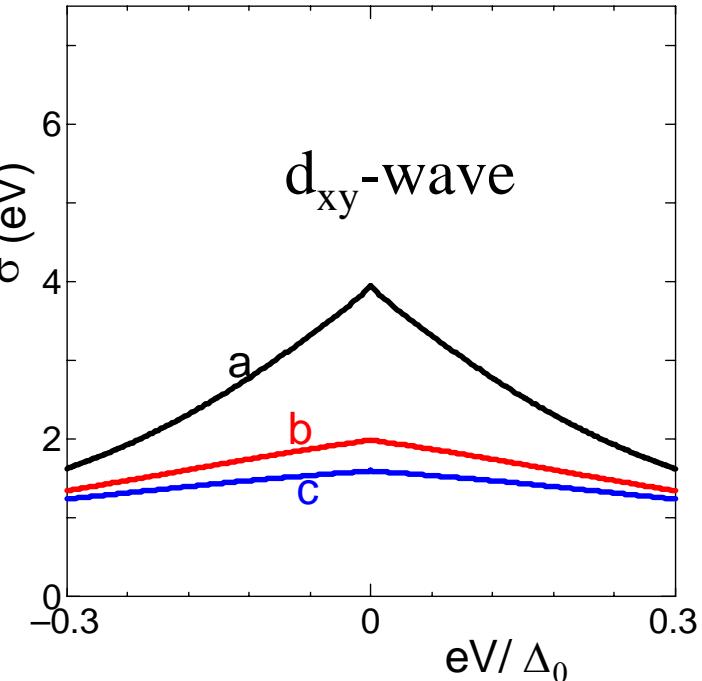


(proximity & MARS)

Normalized tunneling conductance

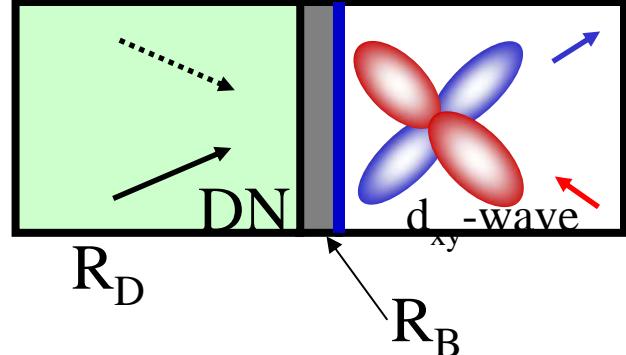


$R_D/R_B=0$
 $R_D/R_B=0.5$
 $R_D/R_B=1$



ZBCP by MARS R_D/R_B is small

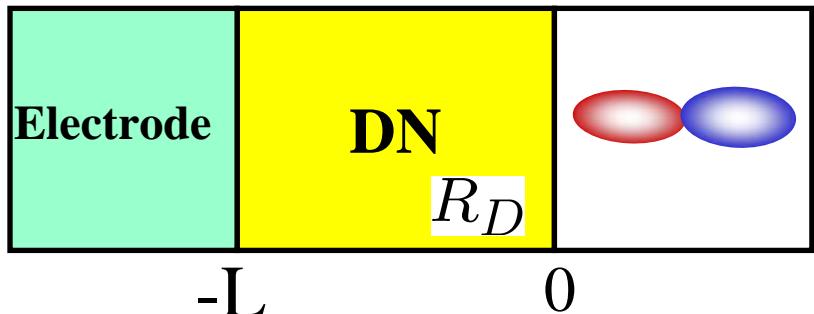
ZBCP by proximity effect
 R_D/R_B is large **Giant ZBCP**



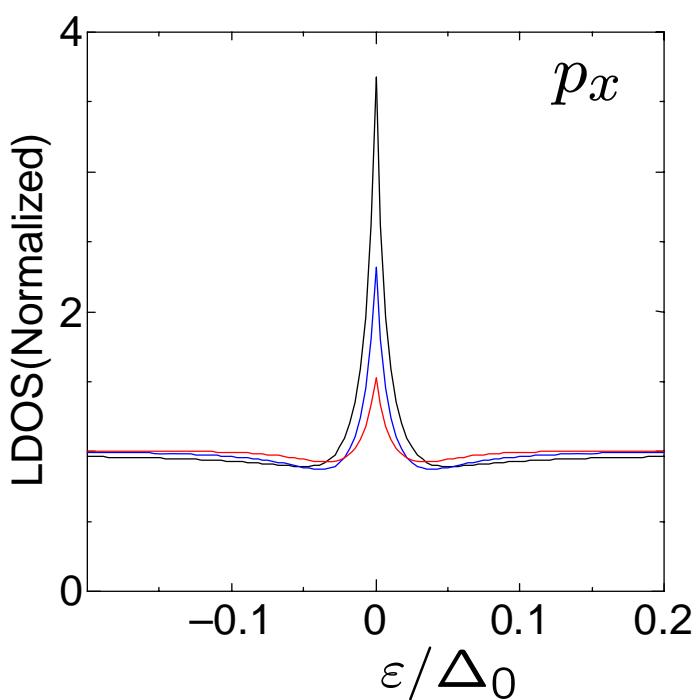
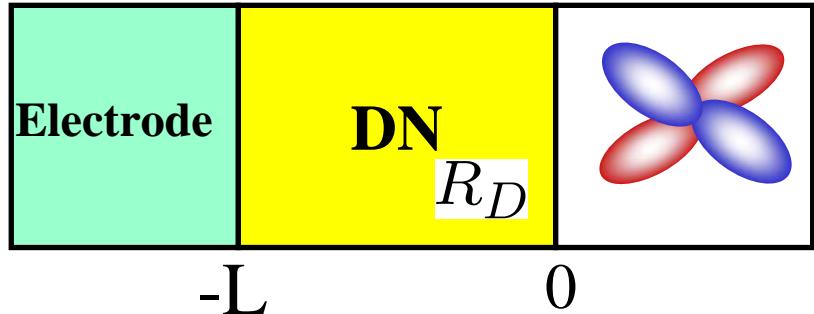
ZBCP only by MARS
No proximity effect

Local density of states in DN

R_B

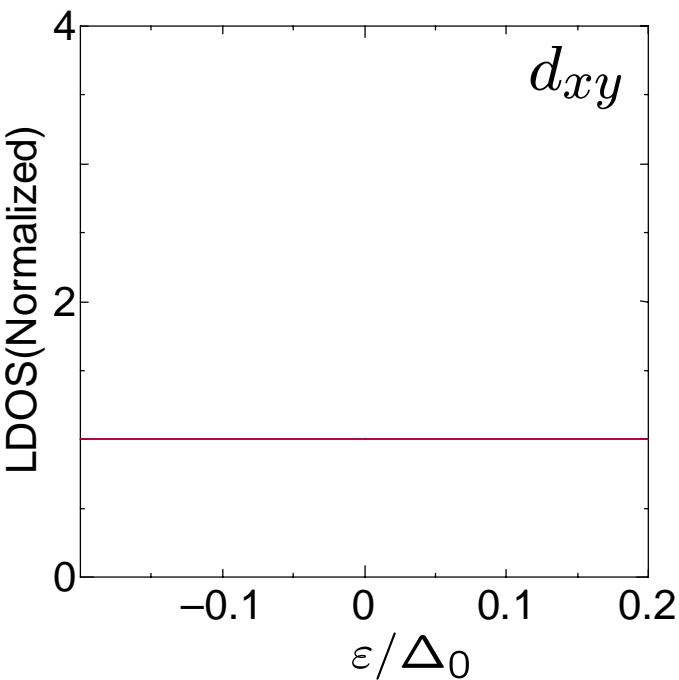


R_B



Z=1.5
 $R_D/R_B = 0.5$
 $E_{th}/\Delta_0 = 0.02$
 $L/\xi = 13$
 $\xi = \sqrt{D/(2\pi T_C)}$

x=0
x=-L/4
x=-L/2



Zero energy peak (ZEP) is expected only for triplet junctions!!

LDOS at $\epsilon = 0$ $\rho(x)$

$$\rho(x) = \cosh\left[\frac{2R_D(x+L)}{LR_0}\right]$$

Unusual energy dependence of pair amplitude

- Triplet junctions

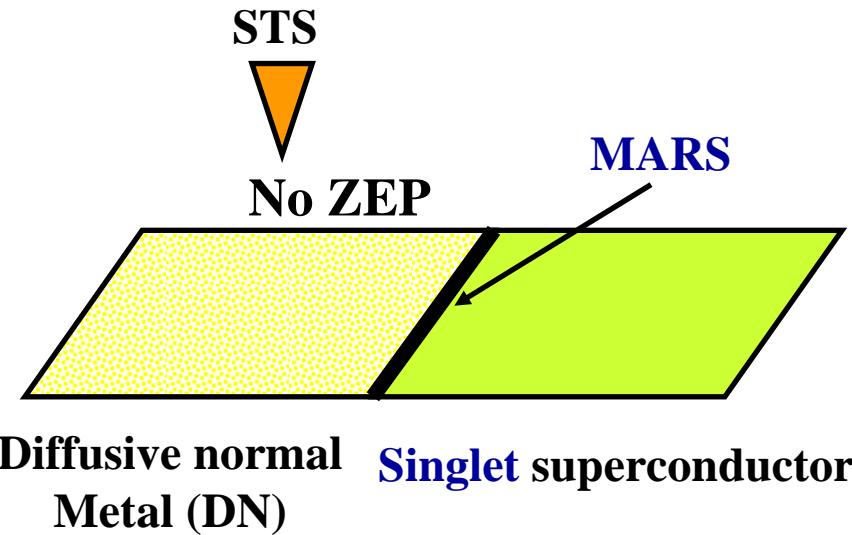
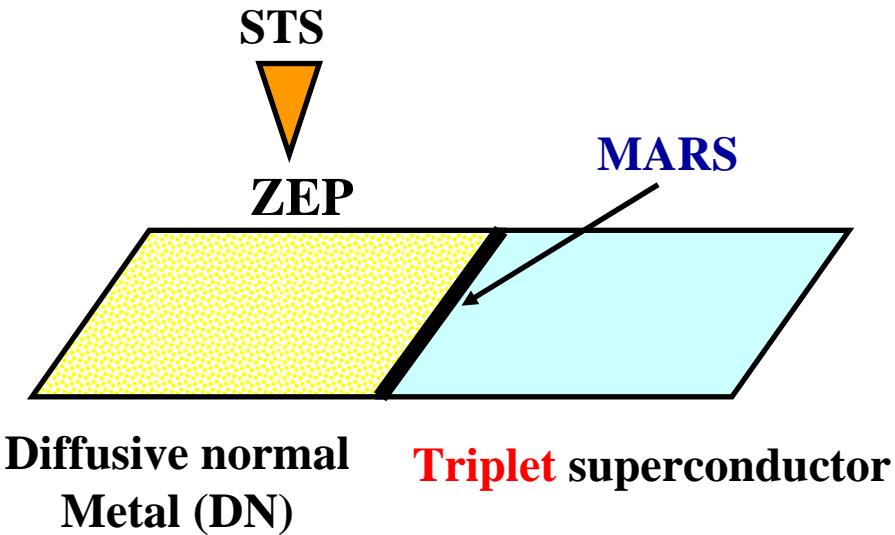
$$f_N(\varepsilon) = -f_N^*(-\varepsilon)$$

- Singlet junctions

$$f_N(\varepsilon) = f_N^*(-\varepsilon)$$

New idea to detect triplet superconductor

MARS (Mid gap Andreev resonance state) can penetrate into DN by **proximity effect** only for triplet superconductor junctions



L DOS in DN has a zero energy peak!!

L DOS in DN does not have a zero energy peak!!

Conclusions

Charge transport in DN/unconventional superconductor junctions

1. Singlet superconductor junctions (Example d-wave)

MARS (Andreev resonant state) competes with proximity effect
MARS can not penetrate into DN

2. Triplet superconductor junctions (Example p-wave)

MARS can coexist with proximity effect
Total resistance of the junction is drastically reduced!!
MARS can penetrate into DN
Enhanced proximity effect is sensitive to applied magnetic field

Future plans

- (1) To clarify anomalous proximity effect
(Meissner effect, unusual energy dependence of pair amplitude) [Phys. Rev. B 71, 094513(2005)]
- (2) Josephson effect
- (3) Ferromagnet junction and spin current
- (4) Noise and full current statistics

- (1) Sr_2RuO_4 -Ru, Sr_2RuO_4 - $\text{Sr}_3\text{Ru}_2\text{O}_7$
- (2) 3K Phase
- (3) STM experiments of Sr_2RuO_4